#### Music of the Microspheres

#### Eigenvalue problems from micro-gyro design

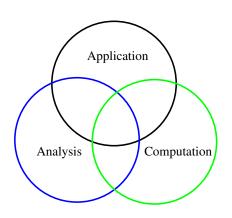
#### David Bindel

Department of Computer Science Cornell University

UTRC, 19 Nov 2013

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### The Computational Science & Engineering Picture

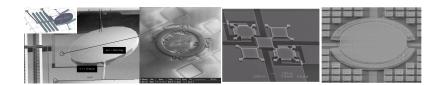


- MEMS (today)
- Smart grids
- Networks

- Linear algebra
- Approximation theory
- Symmetry + structure
- HPC / cloud
- Simulators
- Little languages

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### A Favorite Application: MEMS



#### I've worked on this for a while:

- SUGAR (early 2000s) SPICE for MEMS
- HiQLab (2006) high-Q mechanical resonator device modeling
- AxFEM (2012) solid-wave gyro device modeling

Goal today: an illustrative snapshot.



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### G. H. Bryan (1864–1928)



- Fellow of the Royal Society (1895)
- Stability in Aviation (1911)
- Thermodynamics, hydrodynamics

Bryan was a friendly, kindly, very eccentric individual...

(Obituary Notices of the FRS)

... if he sometimes seemed a colossal buffoon, he himself did not help matters by proclaiming that he did his best work under the influence of alcohol.

(Williams, J.G., The University College of North Wales, 1884–1927)



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### Bryan's Experiment





"On the beats in the vibrations of a revolving cylinder or bell" by G. H. Bryan, 1890

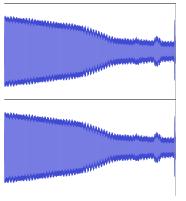
UTRC 5 / 46

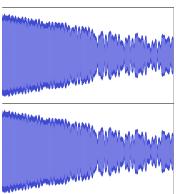
# Bryan's Experiment Today



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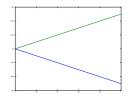
#### The Beat Goes On

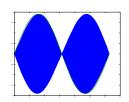




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#### The Beat Goes On





Free vibrations in a rotating frame (simplified):

$$\ddot{\mathbf{q}} + 2\beta\Omega\mathbf{J}\dot{\mathbf{q}} + \omega_0^2\mathbf{q} = 0, \qquad \mathbf{J} \equiv \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}$$

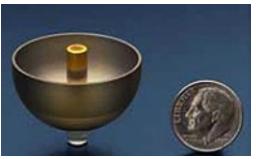
Eigenvalue problem:  $\left(-\omega^2\mathbf{I}+2i\omega\beta\Omega\mathbf{J}+\omega_0^2\right)q=0.$ 

Solutions:  $\omega \approx \Omega_0 \pm \beta \Omega$ .  $\Longrightarrow$  beating  $\propto \Omega!$ 

◆ロ > ◆昼 > ◆差 > ◆差 > を の Q (\*)

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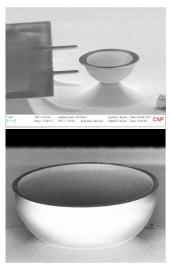
### A Small Application

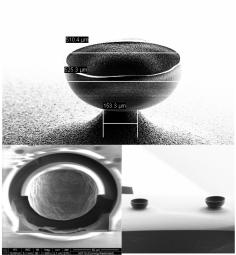


Northrup-Grummond HRG (developed c. 1965–early 1990s)

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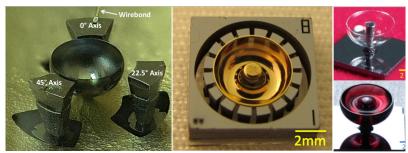
# A Smaller Application (Cornell)

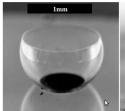


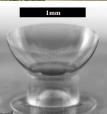


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## A Smaller Application (UMich, GA Tech, Irvine)









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#### Micro-HRG / GOBLiT / OMG





- Goal: Cheap, small (1mm) HRG
- Collaborator roles:
  - Basic design
  - Fabrication
  - Measurement
- Our part:
  - Detailed physics
  - Fast software
  - Sensitivity
  - Design optimization



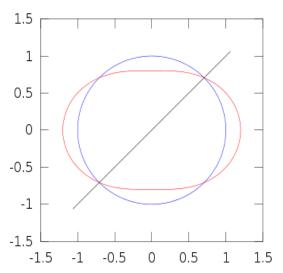
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#### Foucault in Solid State



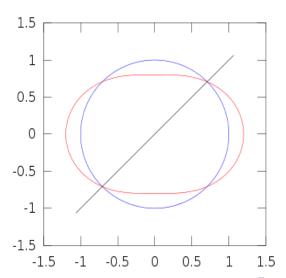
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### Rate Integrating Mode

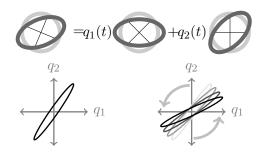


**UTRC** 

### Rate Integrating Mode



#### A General Picture

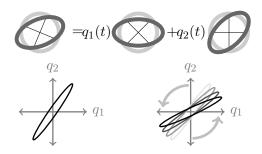


$$\ddot{\mathbf{q}} + 2\beta\Omega\mathbf{J}\dot{\mathbf{q}} + \omega_0^2\mathbf{q} = 0, \qquad \mathbf{J} \equiv \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}$$



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#### A General Picture

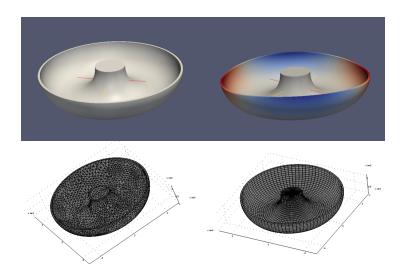


$$\begin{bmatrix} q_1(t) \\ q_2(t) \end{bmatrix} \approx \begin{bmatrix} \cos(-\beta\Omega t) & -\sin(-\beta\Omega t) \\ \sin(-\beta\Omega t) & \cos(-\beta\Omega t) \end{bmatrix} \begin{bmatrix} q_1^0(t) \\ q_2^0(t) \end{bmatrix}.$$



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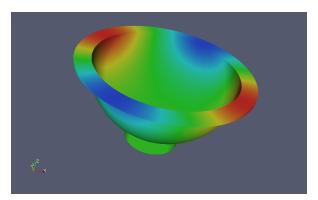
# An Uncritical FEA Approach





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#### The Perturbation Picture

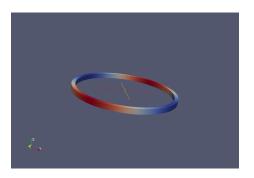


#### Perturbations split degenerate modes:

- Coriolis forces (good)
- Imperfect fab (bad, but physical)
- Discretization error (non-physical)

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### Three Step Program

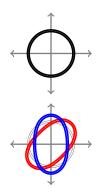


- Perfect geometry, no rotation
- Perfect geometry, rotation
- Imperfect geometry



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# Step I: Perfect Geometry, No Rotation



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### Step I: Perfect Geometry, No Rotation

Free vibration problem in weak form

$$\forall \mathbf{w}, \quad b(\mathbf{w}, \ddot{\mathbf{u}}) + a(\mathbf{w}, \mathbf{u}) = 0.$$

Symmetry: Q any rotation or reflection

$$b(Q\mathbf{w}, Q\mathbf{u}) = b(\mathbf{w}, \mathbf{u})$$

$$a(Q\mathbf{w}, Q\mathbf{u}) = a(\mathbf{w}, \mathbf{u})$$

Decompose by invariant subspaces of  $Q \implies$  Fourier analysis!

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### Fourier Expansion and Axisymmetric Shapes

Decompose into symmetric  $\mathbf{u}^c$  and antisymmetric  $\mathbf{u}^s$  in y:

$$\mathbf{u}^c = \sum_{m=0}^{\infty} \mathbf{\Phi}_m^c(\theta) \mathbf{u}_m^c(r,z), \qquad \quad \mathbf{u}^s = \sum_{m=0}^{\infty} \mathbf{\Phi}_m^s(\theta) \mathbf{u}_m^s(r,z)$$

where

$$\Phi_m^c(\theta) = \operatorname{diag}(\cos(m\theta), \sin(m\theta), \cos(m\theta)) 
\Phi_m^s(\theta) = \operatorname{diag}(-\sin(m\theta), \cos(m\theta), -\sin(m\theta)).$$

Modes involve only one azimuthal number m; degenerate for m > 1.

Preserve structure in FE: shape functions  $N_j(r,z)\Phi_m^{c,s}(\theta)$ 



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### **Block Diagonal Structure**

Finite element system:  $\mathbf{M}\ddot{\mathbf{u}}^h + \mathbf{K}\mathbf{u}^h = 0$ 

Mass has same structure.



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#### Step II: Perfect Geometry, Rotation

Free vibration problem in weak form

$$\forall \mathbf{w}, \quad b(\mathbf{w}, \mathbf{a}) + a(\mathbf{w}, \mathbf{u}) = 0.$$

where

$$\mathbf{a} = \ddot{\mathbf{u}} + 2\mathbf{\Omega} \times \dot{\mathbf{u}} + \mathbf{\Omega} \times (\mathbf{\Omega} \times \mathbf{x}) + \dot{\mathbf{\Omega}} \times \mathbf{x}$$

Discretize by finite elements as before:

$$\mathbf{M}\ddot{\mathbf{u}}^h + \mathbf{C}\dot{\mathbf{u}}^h + \mathbf{K}\ddot{\mathbf{u}}^h = 0$$

where C comes from Coriolis term  $(2b(\mathbf{w}, \mathbf{\Omega} \times \dot{\mathbf{u}}))$ .

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#### Block Structure of Finite Element Matrix

Discretize  $2b(\mathbf{w}, \mathbf{\Omega} \times \dot{\mathbf{u}})$ :

$$\mathbf{C} = \begin{bmatrix} \mathbf{C}_{00} & \mathbf{C}_{01} \\ \mathbf{C}_{10} & \mathbf{C}_{11} & \mathbf{C}_{12} \\ & \mathbf{C}_{21} & \mathbf{C}_{22} & \mathbf{C}_{23} \\ & & \ddots & \ddots & \ddots \end{bmatrix}$$

Off-diagonal blocks come from cross-axis sensitivity:

$$\mathbf{\Omega} = \mathbf{\Omega}_z \mathbf{e}_z + \mathbf{\Omega}_{r\theta} = \begin{bmatrix} 0 \\ 0 \\ \Omega_z \end{bmatrix} + \begin{bmatrix} \cos(\theta)\Omega_x - \sin(\theta)\Omega_y \\ \sin(\theta)\Omega_x + \cos(\theta)\Omega_y \\ 0 \end{bmatrix}.$$

Neglect cross-axis effects  $(O(\Omega^2/\omega_0^2)$ , like centrifugal effect).

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### Analysis in Ideal Case

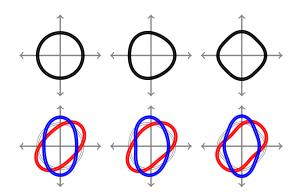
Only need to mesh a 2D cross-section!

- ullet Compute an operating mode  ${f u}_c$  for the non-rotating geometry.
- Compute associated modal mass and stiffness m and k.
- Compute  $g = b(\mathbf{u}_c, e_z \times \mathbf{u}_s)$ .
- ullet Model: motion is approximately  $q_1 \mathbf{u}_c + q_2 \mathbf{u}_s$ , and

$$m\ddot{\mathbf{q}} + 2g\Omega\mathbf{J}\dot{\mathbf{q}} + k\mathbf{q} = 0,$$

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# Step III: Imperfect Geometry



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### What Imperfections?

Let me count the ways...

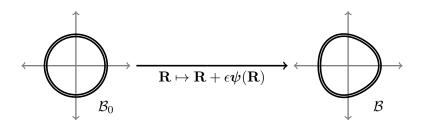
- Over/under etch
- Mask misalignment
- Thickness variations
- Anisotropy of etching single-crystal Si

These are *not* arbitrary!



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### Representing the Perturbation



Write weak form in  $\mathcal{B}_0$  geometry:

$$b(\mathbf{w}, \mathbf{a}) = \int_{\mathcal{B}_0} \rho \mathbf{w} \cdot \mathbf{a} J d\mathcal{B}_0,$$
  
$$a(\mathbf{w}, \mathbf{u}) = \int_{\mathcal{B}_0} \varepsilon(\mathbf{w}) : \mathsf{C} : \varepsilon(\mathbf{u}) J d\mathcal{B}_0,$$

where  $J = \det(\mathbf{I} + \epsilon \mathbf{F})$  with  $\mathbf{F} = \partial \psi / \partial \mathbf{R}$ .



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### Decomposing $\psi$

Do Fourier decomposition of  $\psi$ , too! Consider case where

m = only azimuthal number of  $\mathbf{w}$  n = only azimuthal number of  $\mathbf{u}$ p = only azimuthal number of  $\boldsymbol{\psi}$ 

Then we have selection rules

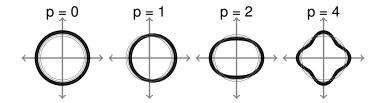
$$a(\mathbf{w}, \mathbf{u}) = \begin{cases} O(\epsilon^k), & |m - n| = kp \\ 0, & \text{otherwise} \end{cases}$$

Similar picture for b.



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# Typical Fabrication Imperfections



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#### Block matrix structure

Ex: 
$$p = 2$$

$$\mathbf{K} = \begin{bmatrix} \mathbf{K}_0 & \epsilon & \epsilon^2 & \epsilon^3 \\ & \mathbf{K}_1 & \epsilon & \epsilon^2 \\ \epsilon & \mathbf{K}_2 & \epsilon & \epsilon^2 \\ & \epsilon & \mathbf{K}_3 & \epsilon \\ \epsilon^2 & \epsilon & \mathbf{K}_4 & \epsilon \\ & \epsilon^2 & \epsilon & \mathbf{K}_5 \\ & \epsilon^3 & \epsilon^2 & \epsilon & \mathbf{K}_6 \\ & & & \ddots \end{bmatrix}$$

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#### Impact of Selection Rules

- Fast FEA: Can neglect most wave numbers / blocks
- Also qualitative information

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#### Qualitative Information

Operating wave number m, perturbation number p:

p=2m	frequencies split by $O(\epsilon)$
kp = 2m	frequencies split at most $O(\epsilon^2)$
p //2m	frequencies change at $O(\epsilon^2)$ , no split
p=1	
$p = 2m \pm 1$	$O(\epsilon)$ cross-axis coupling.

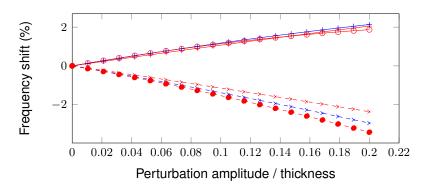
#### Note:

- m=2 affected at first order by p=0 and p=4 (and  $O(\epsilon^2)$  split from p=1 and p=2).
- m=3 affected at first order by p=0 and p=6 (and  $O(\epsilon^2)$  split from p=1 and p=3).



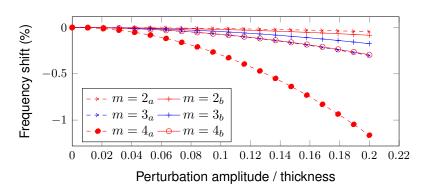
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### Mode Split for Rings: $\psi(r, \theta) = (\cos(2m\theta), 0)$ .



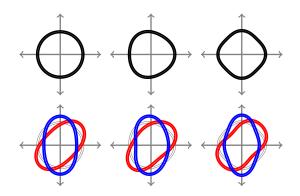
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## Mode Split for Rings: $\psi(r, \theta) = (\cos(m\theta), 0)$ .



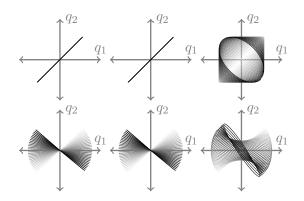
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# **Analyzing Imperfect Rings**



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# **Analyzing Imperfect Rings**



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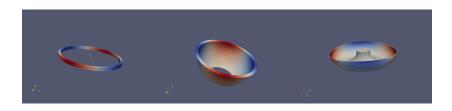
#### Beyond Rings: AxFEM



- Mapped finite / spectral element formulation
- Low-order polynomials through thickness
- High-order polynomials along length
- Trig polynomials in  $\theta$
- Agrees with results reported in literature
- Computes sensitivity to geometry, material parameters, etc.

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#### Read All About It!



Yilmaz and Bindel "Effects of Imperfections on Solid-Wave Gyroscope Dynamics" Proceedings of IEEE Sensors 2013, Nov 3–6.

Thanks to DARPA MRIG + Sunil Bhave and Laura Fegely.



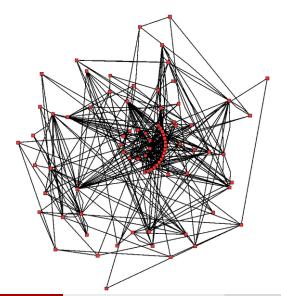
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#### Ad: Declarative HPC and Clouds



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## Ad: Opinions, Game Theory, and Eigenvalues





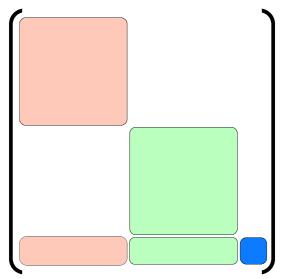
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## Ad: Finding Faults Fast in Smart Grids



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# Ad: Super-Fast PDE Solvers



UTRC 45 / 46

# Ad: Nonlinear Eigenvalues and Resonances





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