

Modeling Resonant Microsystems

Toward Cell Phones on a Chip?

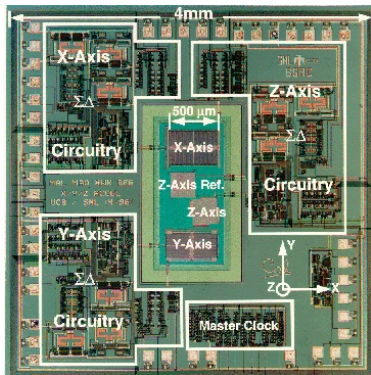
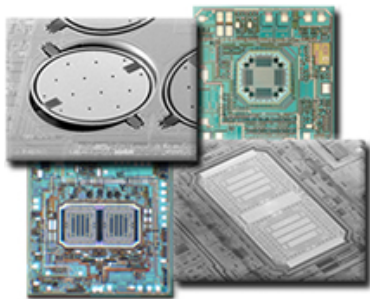
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Department of EECS
University of California, Berkeley

Abel Symposium, 25 May 2006

What are MEMS?

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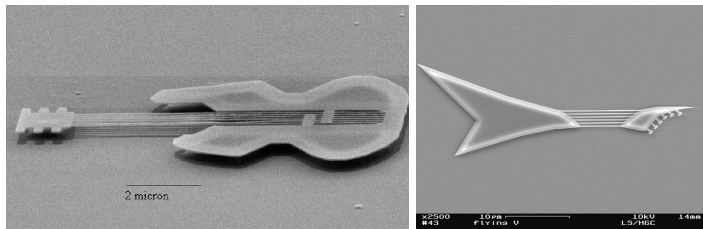
MEMS Basics

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- Micro-Electro-Mechanical Systems
 - Chemical, fluid, thermal, optical (MECFTOMS?)
- Applications:
 - Sensors (inertial, chemical, pressure)
 - Ink jet printers, biolab chips
 - Radio devices: cell phones, inventory tags, pico radio
- Use integrated circuit (IC) fabrication technology
- Tiny, but still classical physics

Resonant MEMS

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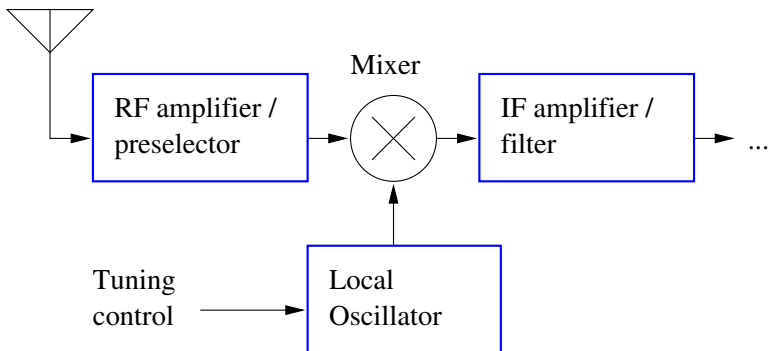


Microguitars from Cornell University (1997 and 2003)

- MHz-GHz mechanical resonators
- Favorite application: radio on chip
- Close second: really high-pitch guitars

The Mechanical Cell Phone

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- Your cell phone has many moving parts!
- What if we replace them with integrated MEMS?

Ultimate Success

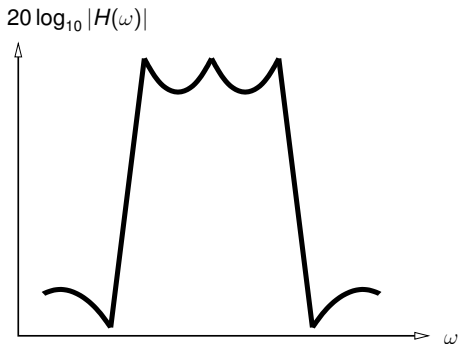
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“Calling Dick Tracy!”



Narrowband Filter Needs

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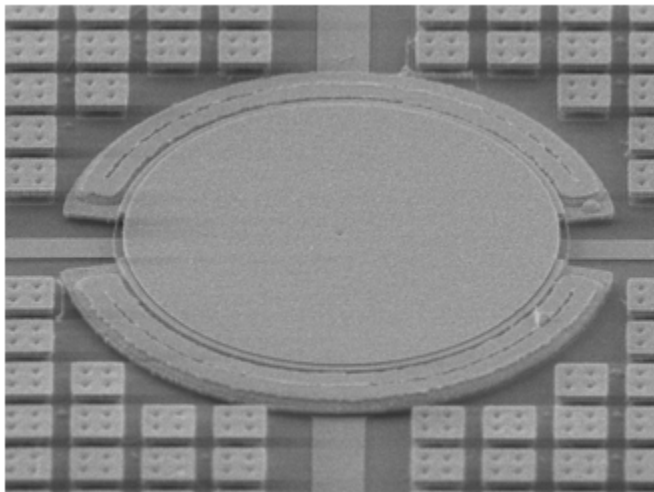


Want building blocks with:

- High frequency
- Low damping
- Tunability

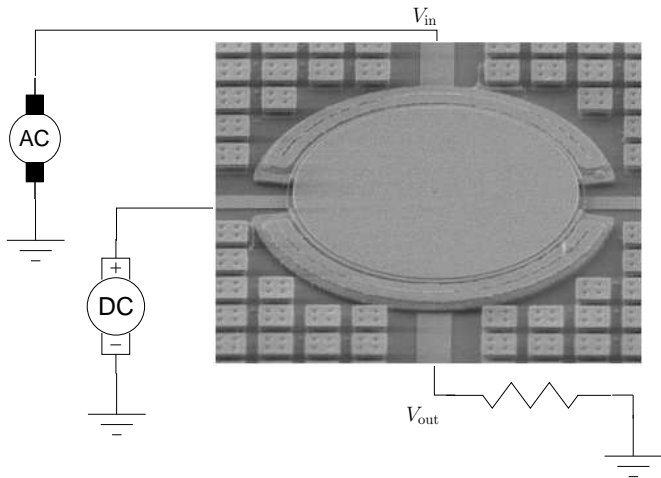
Disk Resonator

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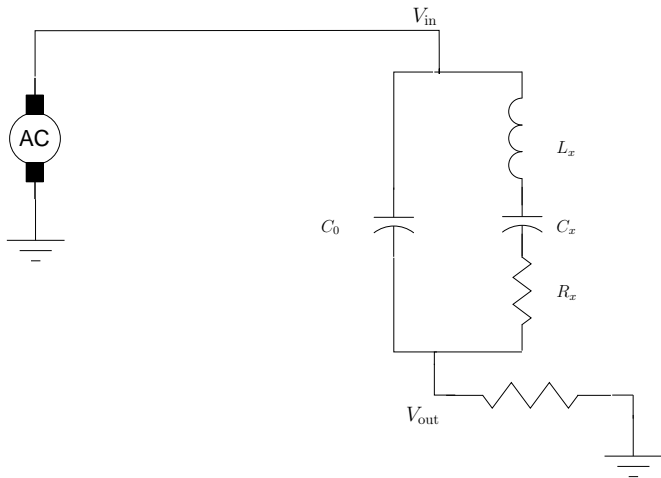
Disk Resonator

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Disk Resonator

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Electromechanical Model

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Kirchoff's current law and balance of linear momentum:

$$\begin{aligned}\frac{d}{dt} (C(u)V) + GV &= I_{\text{external}} \\ Mu_{tt} + Ku - \nabla_u \left(\frac{1}{2} V^* C(u_0) V \right) &= F_{\text{external}}\end{aligned}$$

Linearize about static equilibrium (V_0, u_0) :

$$\begin{aligned}C(u_0) \delta V_t + G \delta V + (\nabla_u C(u_0) \cdot \delta u_t) V_0 &= \delta I_{\text{external}} \\ M \delta u_{tt} + \tilde{K} \delta u + \nabla_u (V_0^* C(u_0) \delta V) &= \delta F_{\text{external}}\end{aligned}$$

where

$$\tilde{K} = K - \frac{1}{2} \frac{\partial^2}{\partial u^2} (V_0^* C(u_0) V_0)$$

Electromechanical Model

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Assume time-harmonic steady state, no external forces:

$$\begin{bmatrix} i\omega C + G & i\omega B \\ -B^T & \tilde{K} - \omega^2 M \end{bmatrix} \begin{bmatrix} \delta \hat{V} \\ \delta \hat{u} \end{bmatrix} = \begin{bmatrix} \delta \hat{l}_{\text{external}} \\ 0 \end{bmatrix}$$

Eliminate the mechanical terms:

$$(i\omega C + G + i\omega B^T (\tilde{K} - \omega^2 M)^{-1} B) \delta \hat{V} = \delta \hat{l}_{\text{external}}$$

Give a name to the coupling transfer function:

$$H(\omega) = B^T (\tilde{K} - \omega^2 M)^{-1} B$$

Goal: Understand electromechanical piece ($i\omega H(\omega)$).

- As a function of geometry and operating point
- Preferably as a simple circuit

Damping and Q

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Designers want high *quality of resonance* (Q)

- Dimensionless damping in a one-dof system

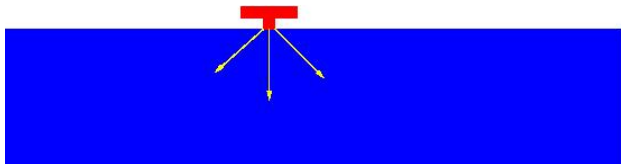
$$\frac{d^2 u}{dt^2} + Q^{-1} \frac{du}{dt} + u = F(t)$$

- For a resonant mode with frequency $\omega \in \mathbb{C}$:

$$Q := \frac{|\omega|}{2 \operatorname{Im}(\omega)} = \frac{\text{Stored energy}}{\text{Energy loss per radian}}$$

Damping Mechanisms

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Possible loss mechanisms:

- Fluid damping
- Material losses
- Thermoelastic damping
- **Anchor loss**

Model substrate as semi-infinite with a

Perfectly Matched Layer (PML).

Perfectly Matched Layers

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- Complex coordinate transformation
- Generates a “perfectly matched” absorbing layer
- Idea works with general linear wave equations
 - Electromagnetics (Bereng er, 1994)
 - Quantum mechanics – *exterior complex scaling* (Simon, 1979)
 - Elasticity in standard finite element framework (Basu and Chopra, 2003)

Model Problem

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- Domain: $x \in [0, \infty)$
- Governing eq:

$$\frac{\partial^2 u}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2} = 0$$

- Fourier transform:

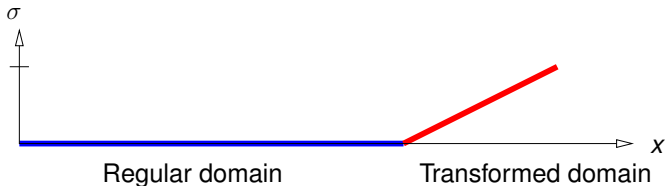
$$\frac{d^2 \hat{u}}{dx^2} + k^2 \hat{u} = 0$$

- Solution:

$$\hat{u} = c_{\text{out}} e^{-ikx} + c_{\text{in}} e^{ikx}$$

Model with Perfectly Matched Layer

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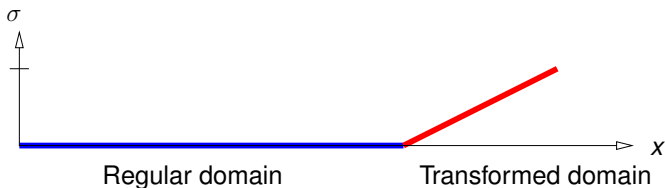
$$\frac{d\tilde{x}}{dx} = \lambda(x) \text{ where } \lambda(s) = 1 - i\sigma(s)$$

$$\frac{d^2\hat{u}}{d\tilde{x}^2} + k^2\hat{u} = 0$$

$$\hat{u} = c_{\text{out}}e^{-ik\tilde{x}} + c_{\text{in}}e^{ik\tilde{x}}$$

Model with Perfectly Matched Layer

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$$\frac{d\tilde{x}}{dx} = \lambda(x) \text{ where } \lambda(s) = 1 - i\sigma(s),$$

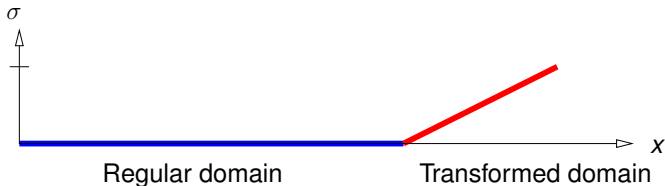
$$\frac{1}{\lambda} \frac{d}{dx} \left(\frac{1}{\lambda} \frac{d\hat{u}}{dx} \right) + k^2 \hat{u} = 0$$

$$\hat{u} = c_{\text{out}} e^{-ikx - k\Sigma(x)} + c_{\text{in}} e^{ikx + k\Sigma(x)}$$

$$\Sigma(x) = \int_0^x \sigma(s) ds$$

Model with Perfectly Matched Layer

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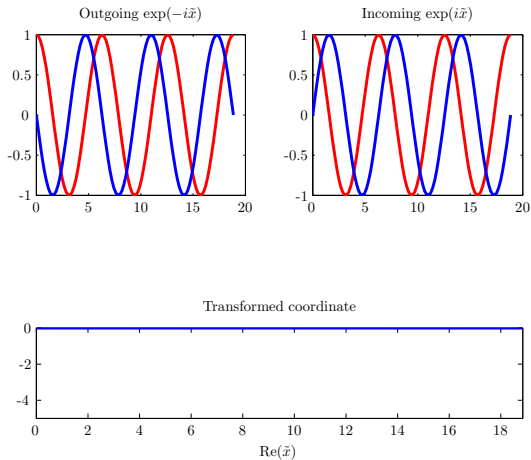


If solution clamped at $x = L$ then

$$\frac{c_{\text{in}}}{c_{\text{out}}} = O(e^{-k\gamma}) \text{ where } \gamma = \Sigma(L) = \int_0^L \sigma(s) ds$$

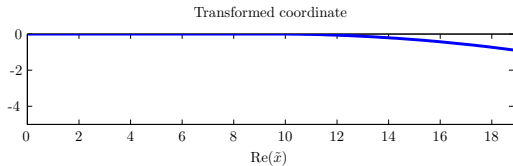
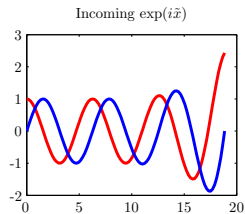
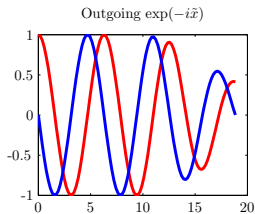
Model Problem Illustrated

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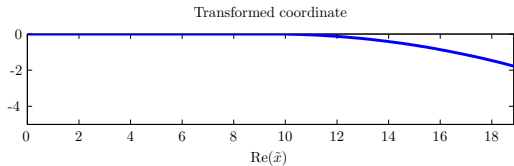
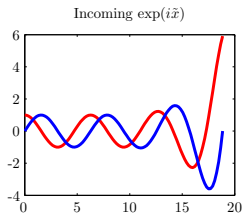
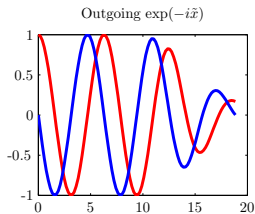
Model Problem Illustrated

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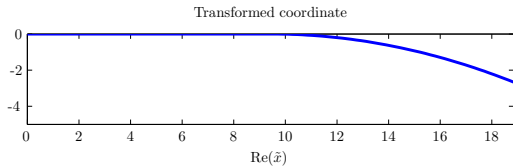
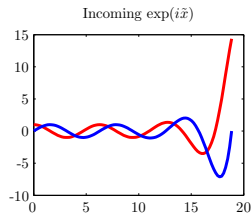
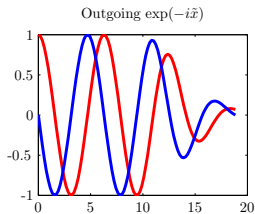
Model Problem Illustrated

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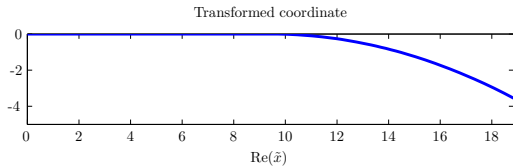
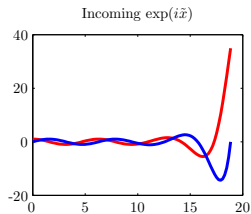
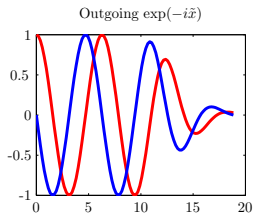
Model Problem Illustrated

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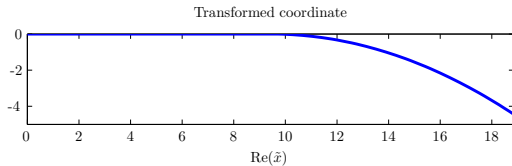
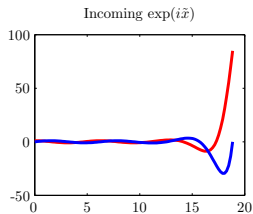
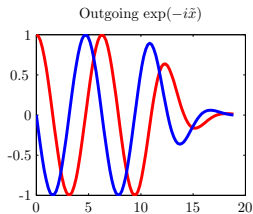
Model Problem Illustrated

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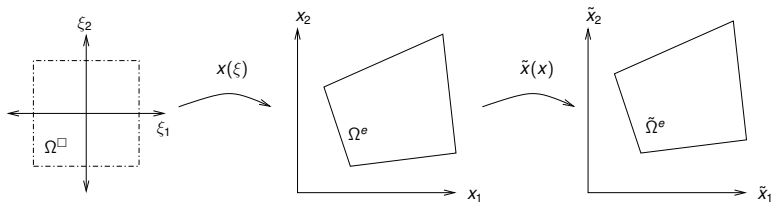
Model Problem Illustrated

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Finite Element Implementation

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- Combine PML and isoparametric mappings

$$\mathbf{k}^e = \int_{\Omega^\square} \tilde{\mathbf{B}}^T \mathbf{D} \tilde{\mathbf{B}} \tilde{J} d\Omega^\square$$

$$\mathbf{m}^e = \int_{\Omega^\square} \rho \mathbf{N}^T \mathbf{N} \tilde{J} d\Omega^\square$$

- Matrices are *complex symmetric*

Eigenvalues and Model Reduction

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Want to know about the transfer function $H(\omega)$:

$$H(\omega) = B^T (K - \omega^2 M)^{-1} B$$

Can either

- Locate poles of H (eigenvalues of (K, M))
- Plot H in a frequency range (Bode plot)

Usual tactic: subspace projection

- Build an Arnoldi basis V
- Compute with much smaller $V^* K V$ and $V^* M V$

Can we do better?

Variational Principles

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- Variational form for complex symmetric eigenproblems:
 - Hermitian (Rayleigh quotient):

$$\rho(v) = \frac{v^* K v}{v^* M v}$$

- Complex symmetric (modified Rayleigh quotient):

$$\theta(v) = \frac{v^T K v}{v^T M v}$$

- First-order accurate eigenvectors \implies
Second-order accurate eigenvalues.
- Key: relation between left and right eigenvectors.

Accurate Model Reduction

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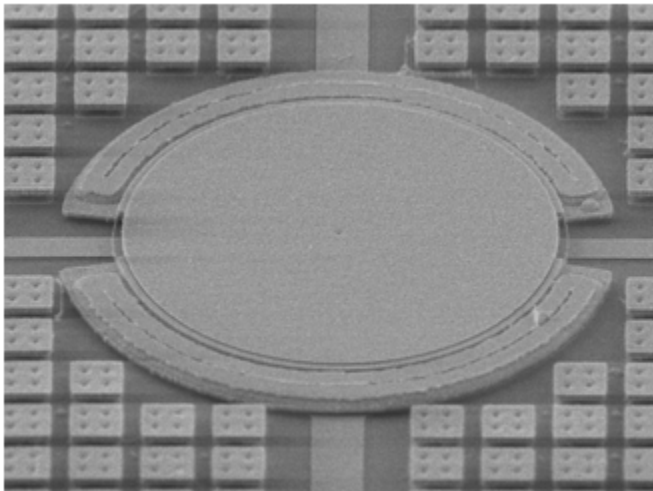
- Build new projection basis from V :

$$W = \text{orth}[\text{Re}(V), \text{Im}(V)]$$

- $\text{span}(W)$ contains both \mathcal{K}_n and $\bar{\mathcal{K}}_n$
 \implies double digits correct vs. projection with V
- W is a real-valued basis
 \implies projected system is complex symmetric

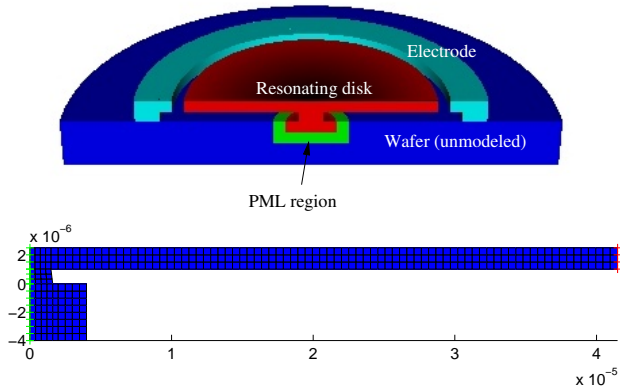
Disk Resonator Simulations

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Disk Resonator Mesh

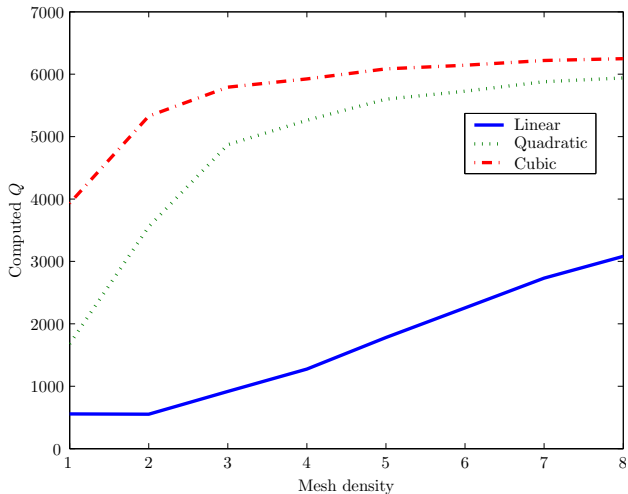
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- Axisymmetric model with bicubic mesh
- About 10K nodal points in converged calculation

Mesh Convergence

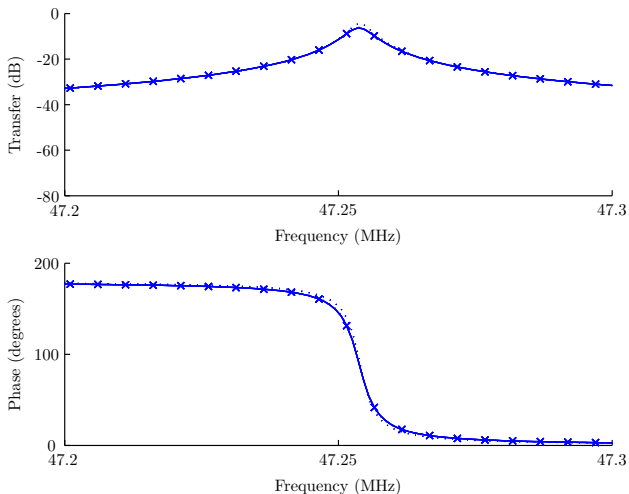
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Cubic elements converge with reasonable mesh density

Model Reduction Accuracy

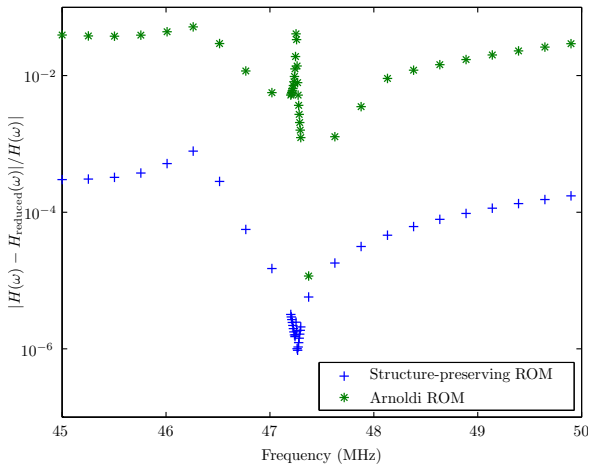
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Results from ROM (solid and dotted lines) nearly indistinguishable from full model (crosses)

Model Reduction Accuracy

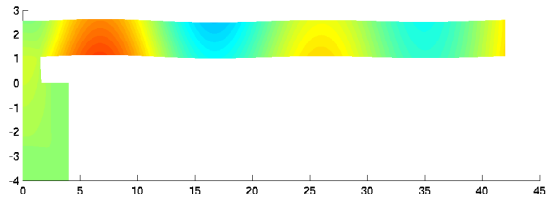
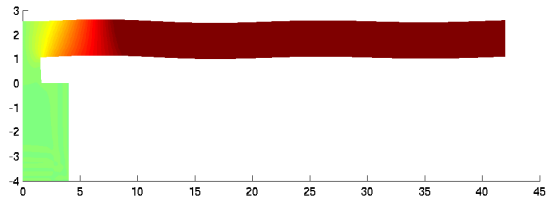
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Preserve structure \implies
get twice the correct digits

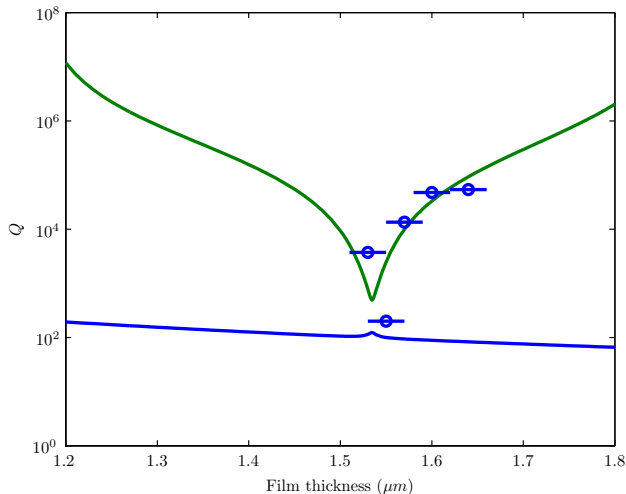
Response of the Disk Resonator

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Variation in Quality of Resonance

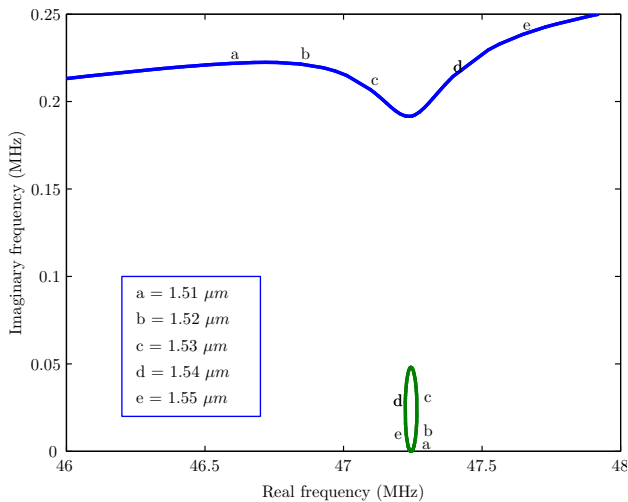
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Simulation and lab measurements vs. disk thickness

Explanation of Q Variation

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Interaction of two nearby eigenmodes

What about:

- Modeling more geometrically complex devices?
- Modeling general dependence on geometry?
- Modeling general dependence on operating point?
- Computing nonlinear dynamics?
- Digesting all this to help designers?

Concluding Thoughts

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The difference between art and science is that science is what we understand well enough to explain to a computer. Art is everything else.

Donald Knuth

The purpose of computing is insight, not numbers.

Richard Hamming