

# Computer Aided Design of Micro-Electro-Mechanical Systems

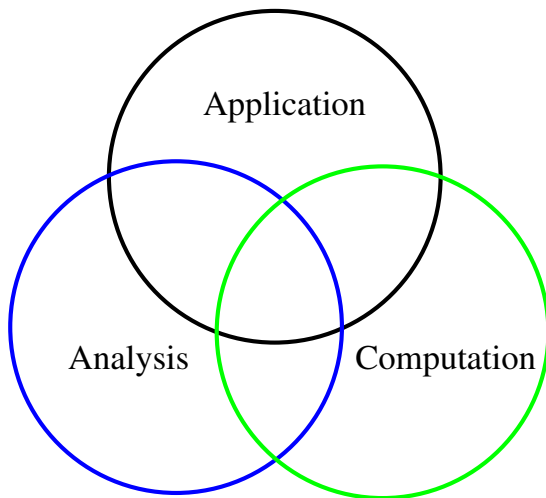
## From Eigenvalues to Devices

David Bindel

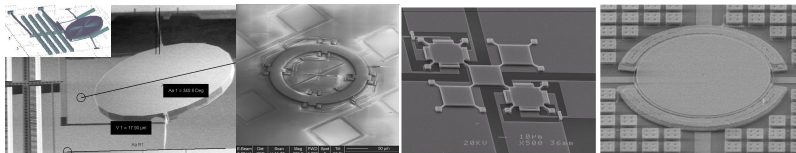
Department of Computer Science  
Cornell University

Duke University, 4 Apr 2013

# The Computational Science & Engineering Picture



# A Favorite Application: MEMS



I've worked on this for a while:

- SUGAR (early 2000s) – SPICE for MEMS
- HiQLab (2006) – high-Q mechanical resonator device modeling
- AxFEM (2012) – solid-wave gyro device modeling

Goal today: two illustrative snapshots.

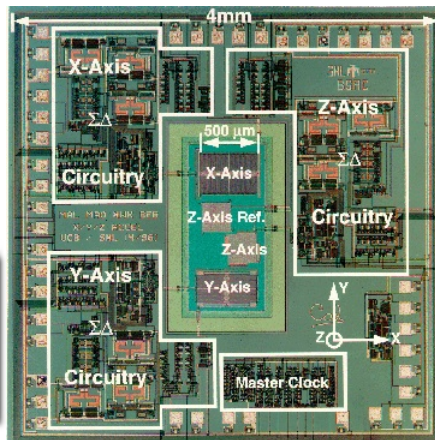
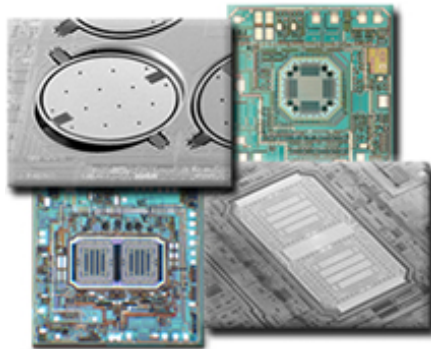
# Outline

- 1 Resonant MEMS
- 2 Anchor losses and disk resonators
- 3 Elastic wave gyros
- 4 Conclusion

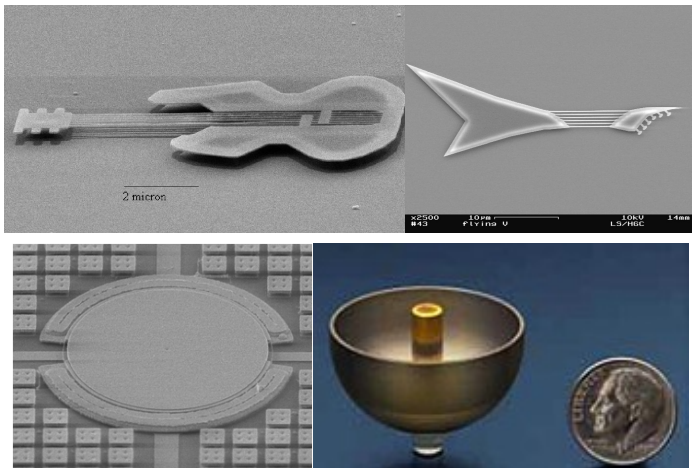
# MEMS Basics

- Micro-Electro-Mechanical Systems
  - Chemical, fluid, thermal, optical (MECFTOMS?)
- Applications:
  - Sensors (inertial, chemical, pressure)
  - Ink jet printers, biolab chips
  - Radio devices: cell phones, inventory tags, pico radio
- Use integrated circuit (IC) fabrication technology
- Tiny, but still classical physics

# Where are MEMS used?



# My favorite applications

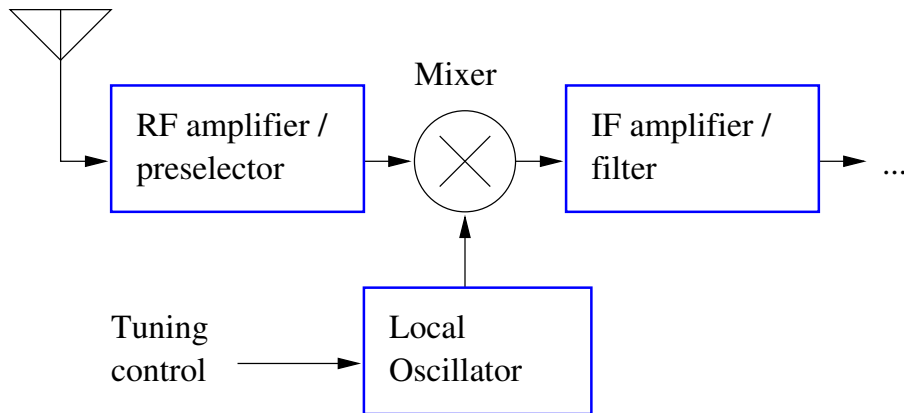


# Why you should care, too!





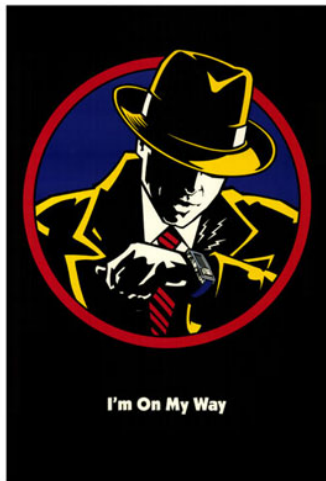
# The Mechanical Cell Phone



- ...and lots of mechanical sensors, too!

# Ultimate Success

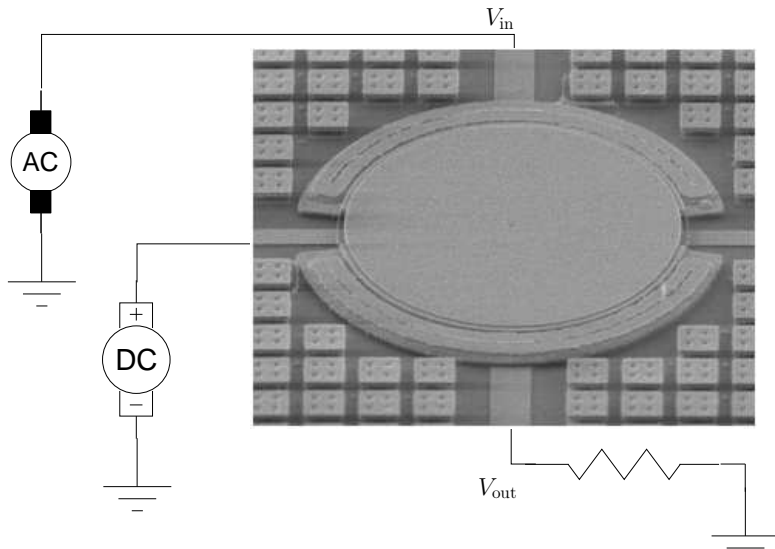
“Calling Dick Tracy!”



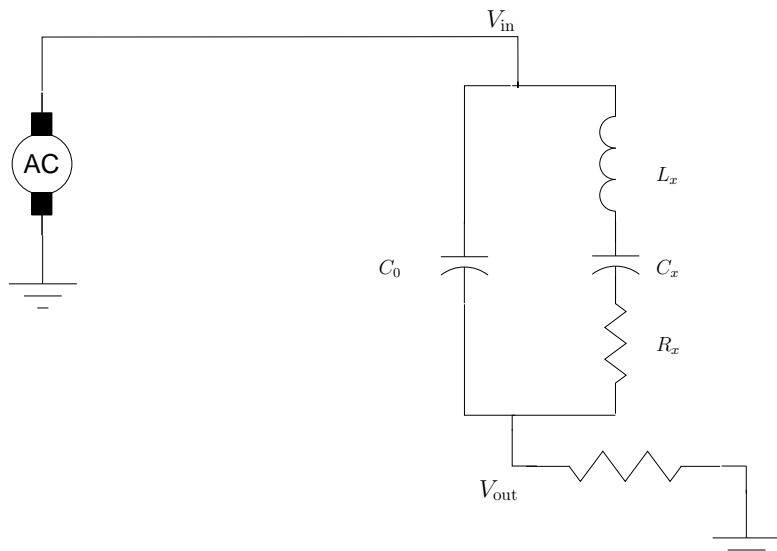
# Computational Challenges

Devices are fun – but I'm not a device designer.  
Why am I in this?

# Model System



# The Circuit Designer View



# Electromechanical Model

Balance laws ( KCL and BLM ):

$$\frac{d}{dt} (C(u)V) + GV = I_{\text{external}}$$

$$Mu_{tt} + Ku - \nabla_u \left( \frac{1}{2} V^* C(u) V \right) = F_{\text{external}}$$

Linearize about static equilibrium  $(V_0, u_0)$ :

$$C(u_0) \delta V_t + G \delta V + (\nabla_u C(u_0) \cdot \delta u_t) V_0 = \delta I_{\text{external}}$$

$$M \delta u_{tt} + \tilde{K} \delta u + \nabla_u (V_0^* C(u_0) \delta V) = \delta F_{\text{external}}$$

where

$$\tilde{K} = K - \frac{1}{2} \frac{\partial^2}{\partial u^2} (V_0^* C(u_0) V_0)$$

# Electromechanical Model

Assume time-harmonic steady state, no external forces:

$$\begin{bmatrix} i\omega C + G & i\omega B \\ -B^T & \tilde{K} - \omega^2 M \end{bmatrix} \begin{bmatrix} \delta \hat{V} \\ \delta \hat{u} \end{bmatrix} = \begin{bmatrix} \delta \hat{I}_{\text{external}} \\ 0 \end{bmatrix}$$

Eliminate the mechanical terms:

$$\begin{aligned} Y(\omega) \delta \hat{V} &= \delta \hat{I}_{\text{external}} \\ Y(\omega) &= i\omega C + G + i\omega H(\omega) \\ H(\omega) &= B^T (\tilde{K} - \omega^2 M)^{-1} B \end{aligned}$$

Goal: Understand electromechanical piece ( $i\omega H(\omega)$ ).

- As a function of geometry and operating point
- Preferably as a simple circuit

# Damping and $Q$

Designers want high *quality of resonance* ( $Q$ )

- Dimensionless damping in a one-dof system

$$\frac{d^2u}{dt^2} + Q^{-1} \frac{du}{dt} + u = F(t)$$

- For a resonant mode with frequency  $\omega \in \mathbb{C}$ :

$$Q := \frac{|\omega|}{2 \operatorname{Im}(\omega)} = \frac{\text{Stored energy}}{\text{Energy loss per radian}}$$

To understand  $Q$ , we need damping models!



# The Designer's Dream

Reality is messy:

- Coupled physics
- ... some poorly understood (damping)
- ... subject to fabrication errors

Ideally, would like:

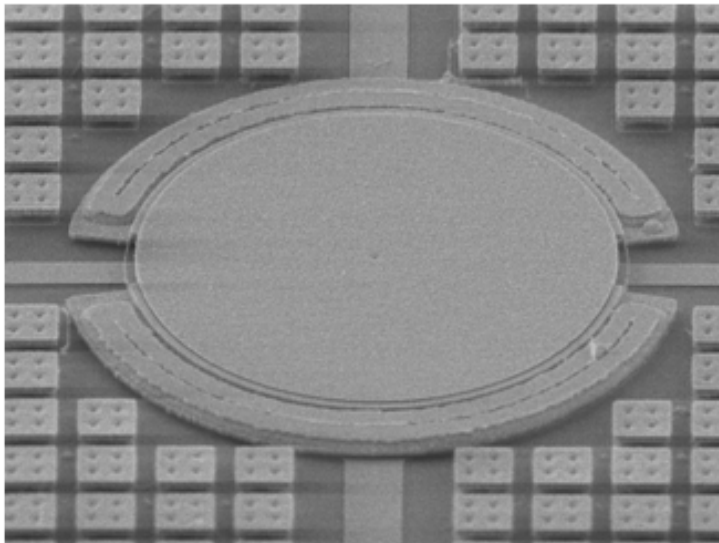
- Simple models for behavioral simulation
- Parameterized for design optimization
- Including all relevant physics
- With reasonably fast and accurate set-up

We aren't there yet.

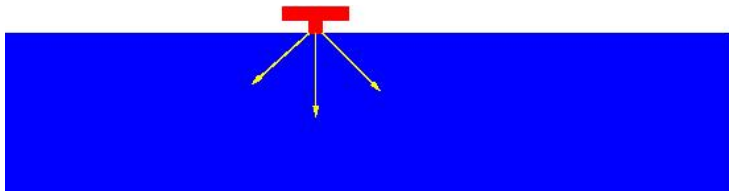
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# Disk Resonator Simulations



# Damping Mechanisms

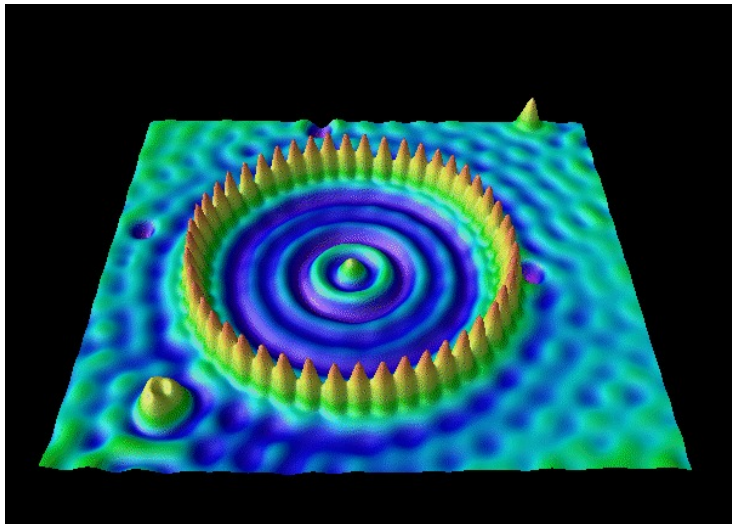


Possible loss mechanisms:

- Fluid damping
- Material losses
- Thermoelastic damping
- **Anchor loss**

Model substrate as semi-infinite  $\implies$  resonances!

# Resonances in Physics



# Resonances and Literature



## Listening to a Monk from Shu Playing the Lute

The monk from Shu with his green lute-case walked  
Westward down Emei Shan, and at the sound  
Of the first notes he strummed for me I heard  
A thousand valleys' rustling pines resound.  
My heart was cleansed, as if in flowing water.  
In bells of frost I heard the resonance die.  
Dusk came unnoticed over the emerald hills  
And autumn clouds layered the darkening sky.

Chinese Poems on the Underground

Li Bai (AD 701-763) translated by Vikram Seth. Three Chinese Poets (Shan) 1992

Calligraphy by Qi Lin Lai

A cultural exchange between Shanghai Metro and London Underground



MAYOR OF LONDON



Transport for London



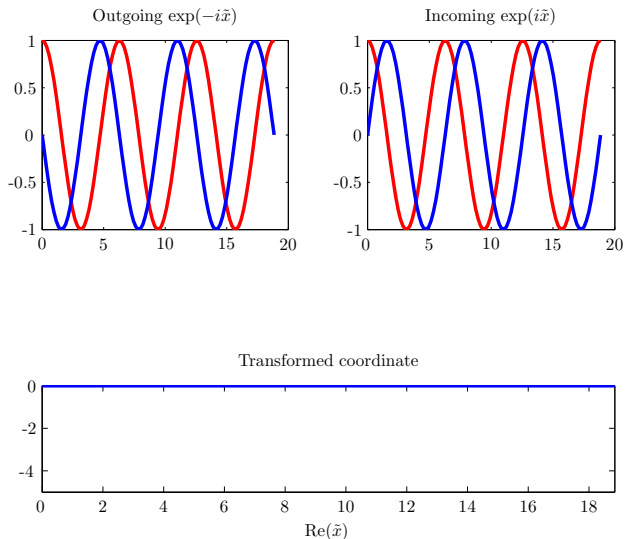
*In bells of frost I heard the resonance die.*

– Li Bai (translated by Vikram Seth)

# Perfectly Matched Layers

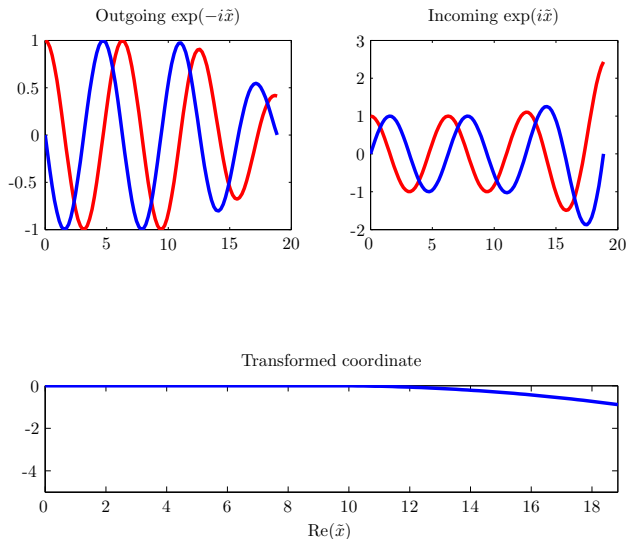
- Complex coordinate transformation
- Generates a “perfectly matched” absorbing layer
- Idea works with general linear wave equations
  - Electromagnetics (Bereng r, 1994)
  - Quantum mechanics – *exterior complex scaling* (Simon, 1979)
  - Elasticity in standard finite element framework (Basu and Chopra, 2003)
  - Works great for MEMS, too! (Bindel and Govindjee, 2005)

# Model Problem Illustrated

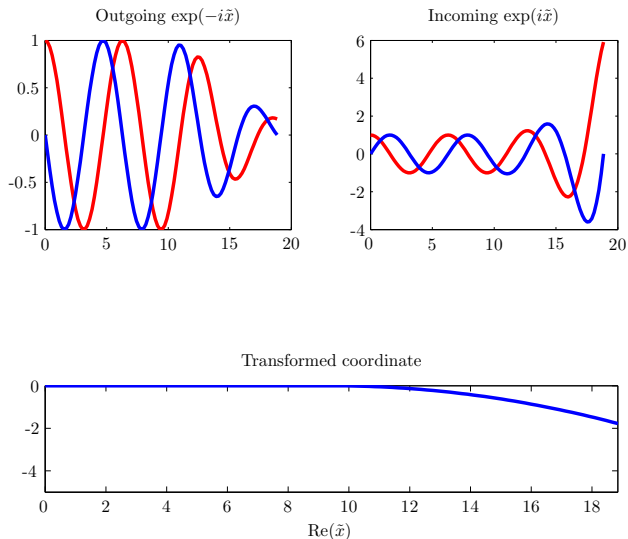




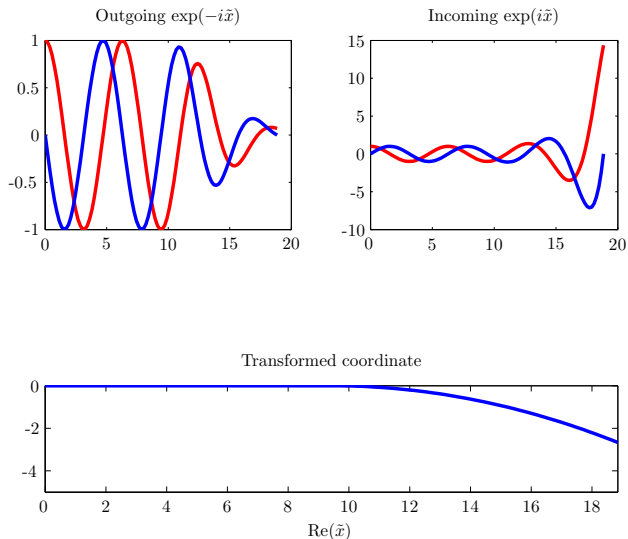
# Model Problem Illustrated



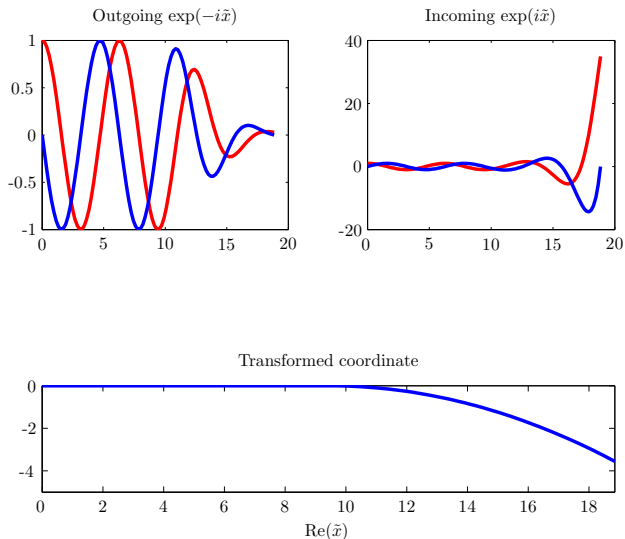
# Model Problem Illustrated



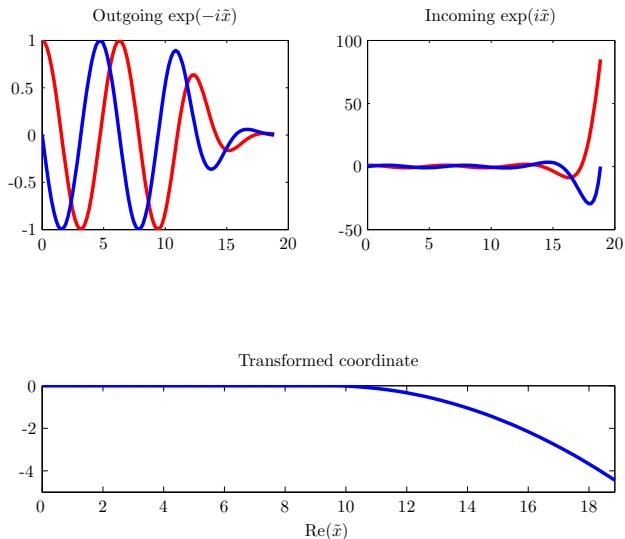
# Model Problem Illustrated



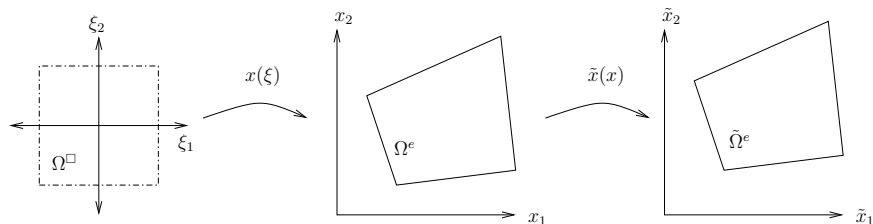
# Model Problem Illustrated



# Model Problem Illustrated

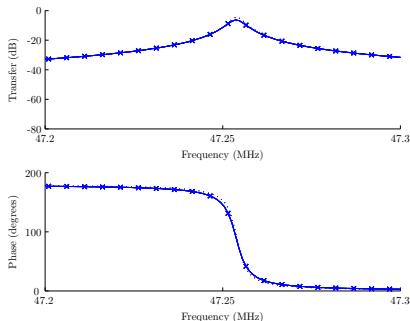


# Finite Element Implementation



Matrices are *complex symmetric*

# Eigenvalues and Model Reduction



Goal: understand  $H(\omega)$ :

$$H(\omega) = B^T(K - \omega^2 M)^{-1} B$$

Look at

- Poles of  $H$  (eigenvalues)
- Bode plots of  $H$

*Model reduction:* Replace  $H(\omega)$  by cheaper  $\hat{H}(\omega)$ .

# Approximation from Subspaces

A general recipe for large-scale numerical approximation:

- 1 A subspace  $\mathcal{V}$  containing good approximations.
- 2 A criterion for “optimal” approximations in  $\mathcal{V}$ .

Basic building block for eigensolvers and model reduction!

Better subspaces, better criteria, better answers.



# Variational Principles

- Variational form for complex symmetric eigenproblems:
  - Hermitian (Rayleigh quotient):

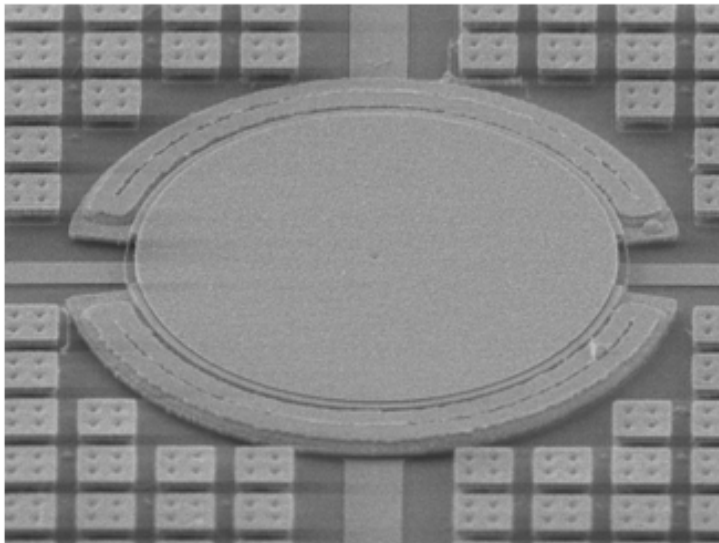
$$\rho(v) = \frac{v^* K v}{v^* M v}$$

- Complex symmetric (modified Rayleigh quotient):

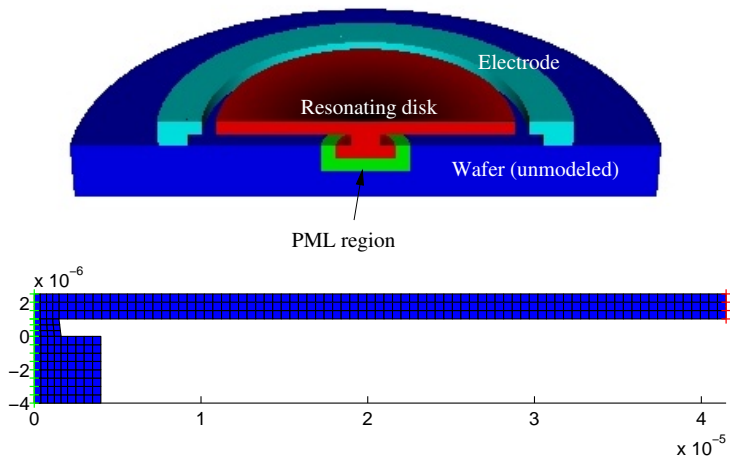
$$\theta(v) = \frac{v^T K v}{v^T M v}$$

- First-order accurate eigenvectors  $\implies$   
Second-order accurate eigenvalues.
- Good for model reduction, too!

# Disk Resonator Simulations

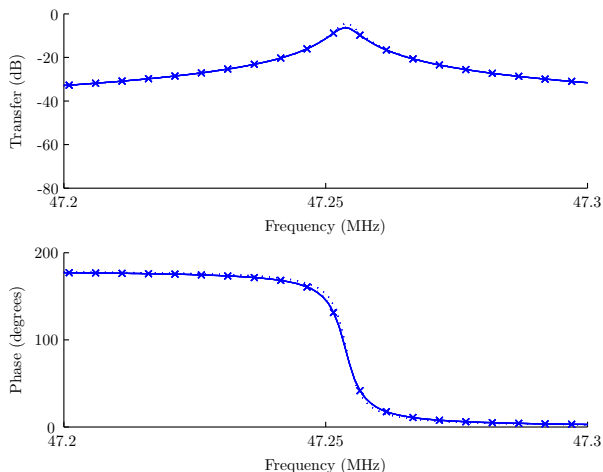


# Disk Resonator Mesh



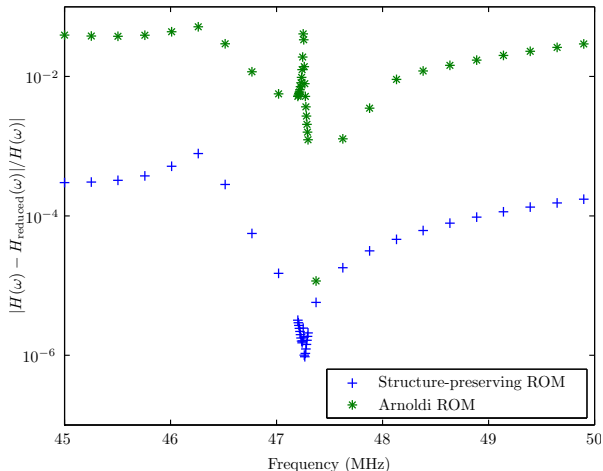
Axisymmetric model, bicubic,  $\approx 10^4$  nodal points at convergence

# Model Reduction Accuracy



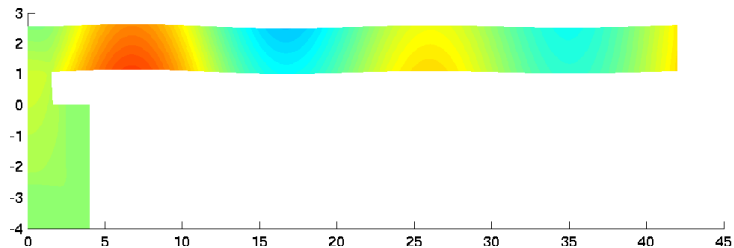
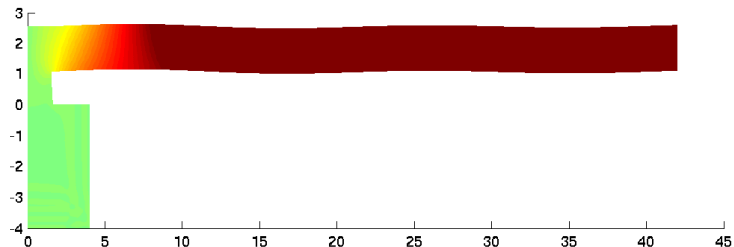
Results from ROM (solid and dotted lines) nearly indistinguishable from full model (crosses)

# Model Reduction Accuracy

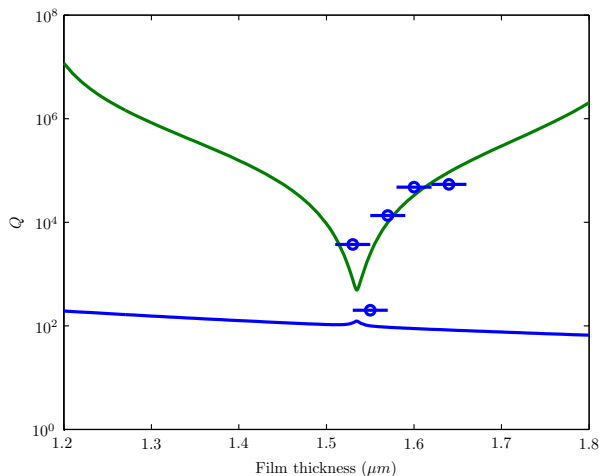


Preserve structure  $\Rightarrow$   
get twice the correct digits

# Response of the Disk Resonator

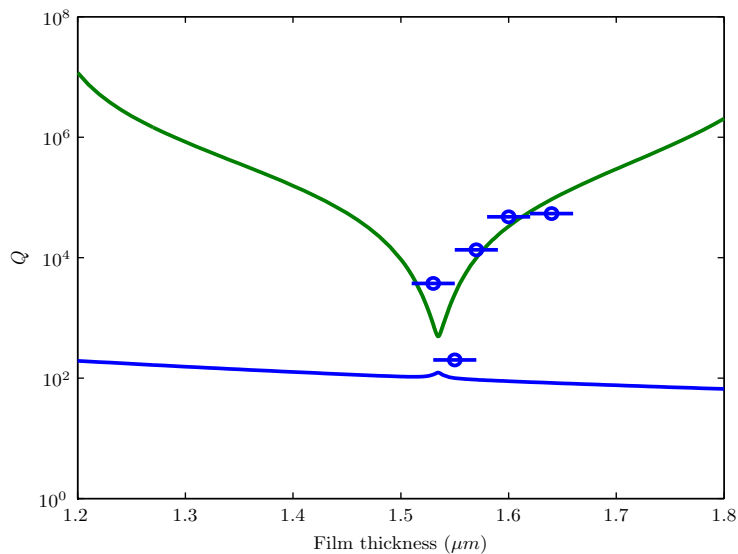


# Variation in Quality of Resonance



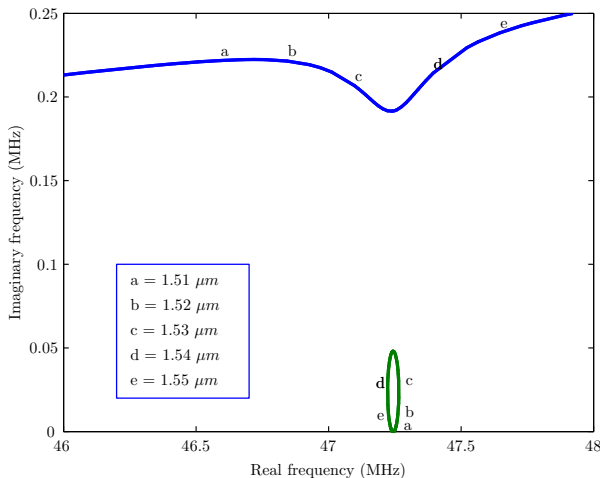
Simulation and lab measurements vs. disk thickness

# Explanation of $Q$ Variation





# Explanation of $Q$ Variation



Interaction of two nearby eigenmodes

# Outline

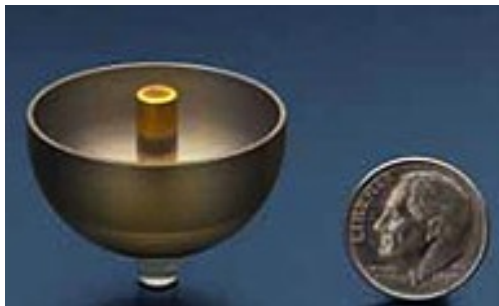
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# Bryan's Experiment



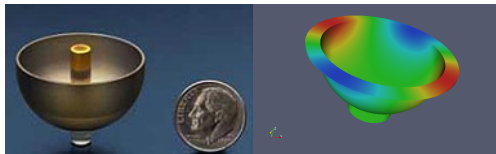
“On the beats in the vibrations of a revolving cylinder or bell”  
by G. H. Bryan, 1890

# A Small Application



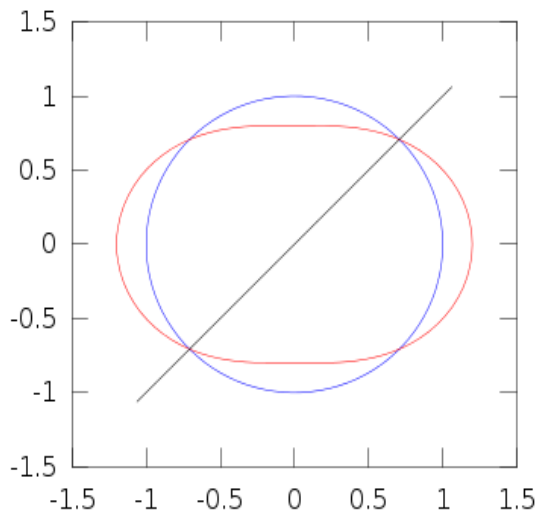
Northrup-Grummond HRG

# Current example: Micro-HRG / GOBLiT / OMG

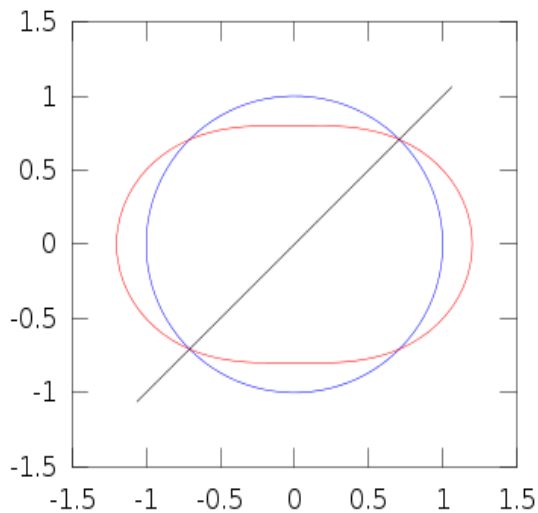


- Goal: Cheap, small (1mm) HRG
- Collaborator roles:
  - Basic design
  - Fabrication
  - Measurement
- Our part:
  - Detailed physics
  - Fast software
  - Sensitivity
  - Design optimization

# How It Works



# How It Works



# Goal state

We want to compute:

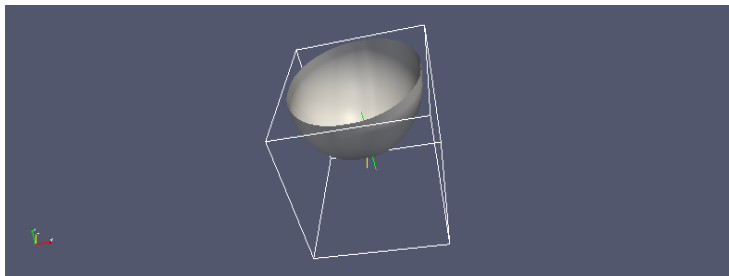
- Geometry
- Fundamental frequencies
- Angular gain (Bryan's factor)
- Damping (thermoelastic, radiation, material)
- Sensitivities of everything
- Effects of symmetry breaking

For speed and accuracy: use structure!

- Axisymmetric geometry  $\implies$  3D to 2D via Fourier
- Perturbed geometry  $\implies$  interactions for different wave numbers

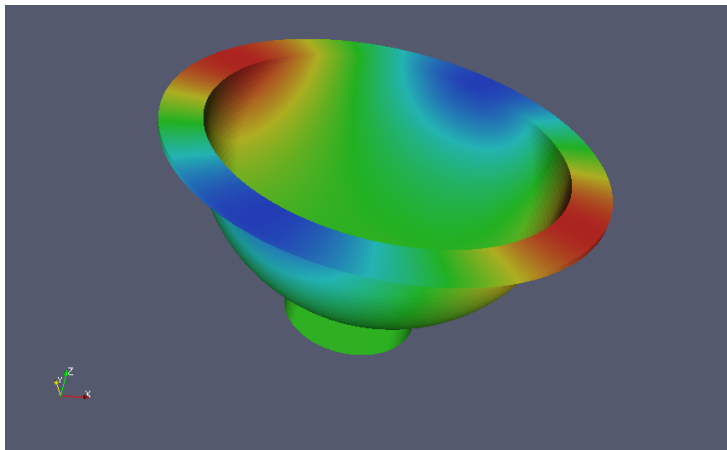


# Getting the Geometry

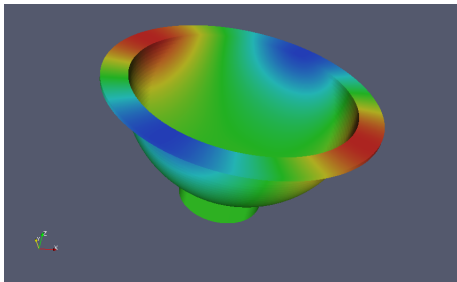


- Simple isotropic etch modeling fails – 1mm is *huge*!
- Working on better simulator (reaction-diffusion).
- For now, take idealized geometries on faith...

# Full Dynamics



# Essential Dynamics



Dynamics in 2D subspace of degenerate modes:

$$(-\omega^2 mI + 2i\omega\Omega gJ + kI) c = 0$$

Scaled gain  $g$  is *Bryan's factor*

$$\text{BF} = \frac{\text{Angular rate of pattern relative to body}}{\text{Angular rate of vibrating body}}$$

*If no parameters in the world were very large or very small,  
science would reduce to an exhaustive list of everything.*  
– Nick Trefethen

# Fourier Picture

Write displacement fields as Fourier series:

$$\mathbf{u} = \sum_{m=0}^{\infty} \left( \begin{bmatrix} u_{mr}(r, z) \cos(m\theta) \\ u_{m\theta}(r, z) \sin(m\theta) \\ u_{mz}(r, z) \cos(m\theta) \end{bmatrix} + \begin{bmatrix} -u'_{mr}(r, z) \sin(m\theta) \\ u'_{m\theta}(r, z) \cos(m\theta) \\ -u'_{mz}(r, z) \sin(m\theta) \end{bmatrix} \right)$$

- Works whenever *geometry* is axisymmetric
- Treat non-axisymmetric geometries as mapped axisymmetric
  - Now *coefficients* in PDEs are non-axisymmetric
- Problems with different  $m$  decouple if *everything* axisymmetric

# Fourier Picture

Perfect axisymmetry:

$$\begin{bmatrix} K_{11} & & & \\ & K_{22} & & \\ & & K_{33} & \\ & & & \ddots \end{bmatrix} - \omega^2 \begin{bmatrix} M_{11} & & & \\ & M_{22} & & \\ & & M_{33} & \\ & & & \ddots \end{bmatrix}$$

# Fourier Picture

Broken symmetry (via coefficients):

$$\begin{bmatrix} K_{11} & \epsilon & \epsilon & \epsilon \\ \epsilon & K_{22} & \epsilon & \epsilon \\ \epsilon & \epsilon & K_{33} & \epsilon \\ \epsilon & \epsilon & \epsilon & \ddots \end{bmatrix} - \omega^2 \begin{bmatrix} M_{11} & \epsilon & \epsilon & \epsilon \\ \epsilon & M_{22} & \epsilon & \epsilon \\ \epsilon & \epsilon & M_{33} & \epsilon \\ \epsilon & \epsilon & \epsilon & \ddots \end{bmatrix}$$

# Perturbing Fourier

Modes “near” azimuthal number  $m$  = nonlinear eigenvalues

$$\left( K_{mm} - \omega^2 M_{mm} + E_{mm}(\omega) \right) u = 0.$$

Need:

- Control on  $E_{mm}$ 
  - Depends on frequency spacing
  - Depends on Fourier analysis of perturbation
- Perturbation theory for nonlinearly perturbed eigenproblems
  - For self-adjoint case, results similar to Lehmann intervals

First-order estimate:  $(K_{mm} - \omega_0^2 M_{mm}) u_0 = 0$ ; then

$$\delta(\omega^2) = \frac{u_0^T E_{mm}(\omega_0) u_0}{u_0^T M_{mm} u_0}.$$



# Perturbation and Radiation

Incorporating numerical radiation BCs gives:

$$\left( K - \omega^2 M + G(\omega) \right) u = 0.$$

Perturbation approach: ignore  $G$  to get  $(\omega_0, u_0)$ . Then

$$\delta(\omega^2) = \frac{u_0^T G(\omega_0) u_0}{u_0^T M_{mm} u_0}.$$

Works when BC has small *influence* (coefficients aren't small).

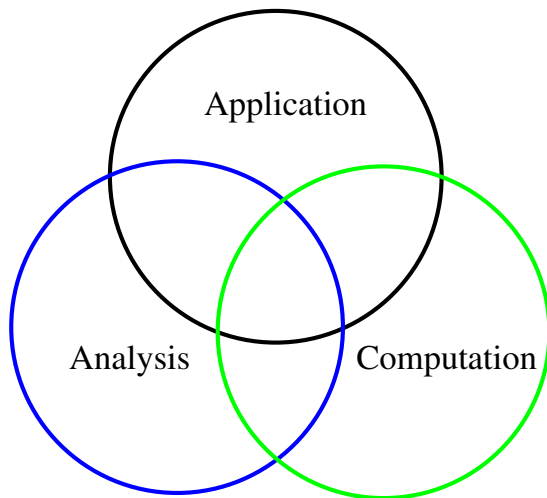
Also an approach to understanding sensitivity to BC!

... explains why PML works okay despite being inappropriate?

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# The Computational Science & Engineering Picture



# Conclusions

*The difference between art and science is that science is what we understand well enough to explain to a computer. Art is everything else.*

*Donald Knuth*

*The purpose of computing is insight, not numbers.*

*Richard Hamming*

- 
- Collaborators:
    - Disk: S. Govindjee, T. Koyama, S. Bhawe, E. Quevy
    - HRG: S. Bhawe, L. Fegely, E. Yilmaz
  - Funding: DARPA MTO, Sloan Foundation