

Automated Phase Mapping with AgileFD and its Application to Light Absorber Discovery in the V–Mn–Nb Oxide System

Santosh K. Suram,[†] Yexiang Xue,[‡] Junwen Bai,[§] Ronan Le Bras,^{‡,⊥} Brendan Rappazzo,[‡] Richard Bernstein,[‡] Johan Bjorck,[‡] Lan Zhou,[†] R. Bruce van Dover,^{||} Carla P. Gomes,^{*,‡} and John M. Gregoire^{*,†,Ⓛ}

[†]Joint Center for Artificial Photosynthesis, California Institute of Technology, Pasadena California 91125, United States

[‡]Department of Computer Science, Cornell University, Ithaca, New York 14850, United States

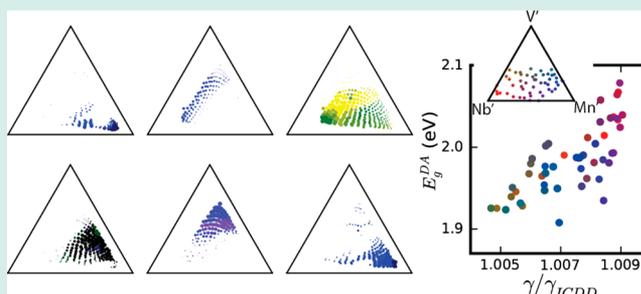
[§]Zhiyuan College, Shanghai Jiao Tong University, Shanghai, China

^{||}Department of Materials Science and Engineering, Cornell University, Ithaca, New York 14850, United States

Supporting Information

ABSTRACT: Rapid construction of phase diagrams is a central tenet of combinatorial materials science with accelerated materials discovery efforts often hampered by challenges in interpreting combinatorial X-ray diffraction data sets, which we address by developing AgileFD, an artificial intelligence algorithm that enables rapid phase mapping from a combinatorial library of X-ray diffraction patterns. AgileFD models alloying-based peak shifting through a novel expansion of convolutional nonnegative matrix factorization, which not only improves the identification of constituent phases but also maps their concentration and lattice parameter as a function of composition. By incorporating Gibbs' phase rule into the algorithm, physically meaningful phase maps are obtained with unsupervised operation, and more refined solutions are attained by injecting expert knowledge of the system. The algorithm is demonstrated through investigation of the V–Mn–Nb oxide system where decomposition of eight oxide phases, including two with substantial alloying, provides the first phase map for this pseudoternary system. This phase map enables interpretation of high-throughput band gap data, leading to the discovery of new solar light absorbers and the alloying-based tuning of the direct-allowed band gap energy of MnV_2O_6 . The open-source family of AgileFD algorithms can be implemented into a broad range of high throughput workflows to accelerate materials discovery.

KEYWORDS: high-throughput screening, machine learning, X-ray diffraction, combinatorial science, band gap tuning



INTRODUCTION

Combinatorial materials science encompasses a suite of techniques for accelerated discovery of functional materials and often entails the establishment of relationships between crystal structure and materials properties.¹ As a result, high-throughput phase mapping, the generation of composition maps of phase concentrations that serve as a proxy for the solid state phase behavior and ultimately the phase diagram for a given synthesis condition, is an important tool for materials discovery, prompting the recent development of synchrotron-based X-ray diffraction (XRD) techniques to acquire 10^2 – 10^4 XRD patterns per day over a composition spread library.² More generally, the accelerated characterization of phase diagrams in high-order composition spaces is a foundational problem in materials science given that many binary phase diagrams are known but few ternary or higher-order phase diagrams have been explored due to the requisite number of experiments. Emboldened by the progress in previous decades in automated XRD analysis for small molecules,³ materials scientists have adopted a variety of techniques in the past 10 years to

automatically generate phase maps (phase concentrations as a function of composition) from XRD patterns, as recently summarized in a perspective article⁴ noting the critical need for improved algorithms that appropriately capture the complexity of materials phase behavior. By developing algorithms at the forefront of artificial intelligence research, we have established a phase mapping tool that models alloying-based shifting of peaks in XRD patterns and incorporates physical constraints based on Gibbs' phase rule. The algorithm performs factor decomposition with an elegant agility that allows users to constrain the spectral demixing based on expert knowledge of the phase behavior, yielding the powerful phase-mapping tool we hereby introduce as AgileFD.

The phase mapping tools developed to date employ a variety of intuitive approaches to the problem of formulating phase maps from a collection of XRD patterns, which generally

Received: October 6, 2016

Revised: November 15, 2016

Published: November 21, 2016

contain a mixture of phase-pure patterns. Recent efforts have demonstrated that if an expert has sufficient knowledge of the phase map to generate training data, in particular expert-labeled XRD patterns for a variety of phases in the system, the trained machine learning model can map phase behavior.⁵ An interesting feature of these supervised machine learning approaches is that phase classification is performed on each XRD pattern sequentially, whereas the remainder of the algorithms discussed below act on the ensemble of XRD patterns to demix multiphase patterns through simultaneous consideration of a wide range of phase mixtures. Clustering XRD patterns based on their similarity⁶ is a powerful data reduction tool but since XRD patterns may be clustered due to either a common phase or a common phase mixture (i.e., when a set of composition samples are from the same phase field in the underlying phase diagram), the behavior of pure-clustering techniques is sensitive to the distance metric and requires delicate manual tuning. Initial efforts to address these issues by using constraint programming (CP) to inject physics-based reasoning into clustering have shown promise but have yet to be developed into an effective phase mapping algorithm for experimental data.⁷ More recently, by performing clustering in a feature vector space that includes the XRD patterns and sample compositions, compositional clusters can be generated that more closely emulate phase diagrams, and this clustering can be performed quickly enough to perform on-the-fly analysis.⁸

An expert's manual analysis of XRD patterns is most closely emulated with a reasoning-based algorithm in which constituent phases are identified by recognizing sets of peaks that coexist in connected composition regions. By expressing the phase mapping problem as a set of logical requirements within a satisfiability modulo theory (SMT) reasoning framework, we previously demonstrated automated phase mapping using pristine, synthetic data.⁹ The implementation of this approach for experimental data suffers from practical limitations due to the computational expense for large data sets and the reliance on accurate peak detection in every XRD pattern. As a result, the development of scalable algorithms for (noisy) experimental data has focused on factoring a collection of XRD patterns into a small set of basis patterns (intended to represent individual phases) and the weighting coefficients or "activation" of the basis patterns for each experimental XRD pattern (intended to represent the amount of each phase present in the corresponding material). Since both XRD patterns and phase concentrations are nonnegative, the mathematical description of this approach is non-negative matrix factorization (NMF) in which the XRD patterns for the composition samples are assembled into a matrix A and the phase map solution corresponds to identifying matrices W (whose columns are the basis patterns) and H (which contains the activation of each basis pattern for each sample), such that

$$A \approx W \cdot H \quad (1)$$

Recent progress in solving the phase mapping problem has thus included implementations of NMF that yield excellent reconstruction of the XRD data set and produce phase maps that adhere (as closely as possible) to the underlying materials physics.

The various NMF approaches can be succinctly summarized by considering their modeling of 3 principle properties of solid-state phase diagrams, particularly for composition libraries synthesized at a fixed temperature: (1) Gibbs' Rule, in

thermodynamic equilibrium, the number of coexisting phases in a given material can be no more than the number of components (number of cations in the case of metal oxides) it contains; (2) connectivity, the set of composition samples containing a given phase must be connected in composition space; and (3) peak shifting or alloying, the interstitial or substitutional solution of elements within a single phase can yield composition-dependent lattice constants that cause XRD patterns to shift as a function of composition. While alloying-based peak shifting can be quite complex for noncubic phases, in the present work we only consider isotropic lattice expansions that are manifested as uniform shifting of peaks in XRD patterns.

Long et al.¹⁰ introduced NMF as a phase mapping strategy, and while this approach is computationally efficient and effective for some data sets, it does not model any of the above properties of phase diagrams, resulting in routine production of nonphysical and uninterpretable phase maps. The complement to bare NMF is CombiFD,¹¹ a powerful factor decomposition approach in which the matrix factorization can be performed under constraints of arbitrary complexity, including the encoding of the above phase diagram properties as hard combinatorial constraints. This approach is analogous to the logic of the SMT algorithm described above and is too computationally intensive for large data sets. GRENDEL¹² is an extension of NMF that was recently adapted for phase mapping by Kusne et al.¹³ to enforce the Connectivity constraint by combining NMF with a graphical model of the composition space. GRENDEL remains computationally efficient by iterating between NMF and clustering with an objective function that combines the loss functions of these subalgorithms, which requires tuning of the relative weights of the loss functions to optimize clustering. Further development of the algorithm is required to enforce Gibbs' Rule, and the reliance on NMF results in an inability to incorporate Peak Shifting. Indeed Peak Shifting has remained a primary challenge in phase mapping, and while powerful approaches for modeling Peak Shifting using dynamic time warping techniques have been proposed,^{7,14} they have yet to be incorporated into scalable phase mapping algorithms.

The AgileFD phase mapping algorithm is based on convolutional non-negative matrix factorization (cNMF),¹⁵ an algorithm that is effective for demixing audio signals, which we have tailored to explicitly incorporate Peak Shifting in the matrix factorization process. The key concept is to introduce multiple copies of the basis pattern for each phase, representing multiple alloys of the same phase but with different lattice constants. Gibbs' Rule is also incorporated in AgileFD-Gibbs, an extension of AgileFD that includes constraint optimization of H using mixed integer programming (MIP). AgileFD also enables an expert user to inject knowledge of the phase map through "expert constraints". We have yet to develop an explicit Connectivity constraint since the phase maps we have generated to date do not substantially violate this rule due to the efficacy of our Gibbs' Rule enforcement. As discussed below, our experience has been that these minor Connectivity violations can be easily corrected by incorporating expert feedback into the solutions.

Perhaps the most powerful aspect of AgileFD is its ability to generate both unsupervised and expert-constrained solutions for sizable data sets (i.e., the data matrix A containing $\sim 10^6$ values corresponding to 100 s of XRD patterns each containing 1000 s of data points) in a few minutes, enabling the user to

interact with solutions via novel data visualizers by altering the number of basis patterns and applying expert constraints until a satisfactory phase map is obtained. The typical analysis of a given XRD data set proceeds by running AgileFD in unsupervised mode to generate candidate phase mappings that are further tailored with constraints applied by expert analysts.

We demonstrate this process flow for a pseudoternary V–Mn–Nb oxide library in which six primary phases are automatically identified with basis patterns that closely match those of known phases. Manual identification of 2 additional, minor phases yields a slightly more refined phase map that is used to interpret high-throughput optical spectroscopy data. While solving the phase behavior of these previously unexplored compositions demonstrates the power of AgileFD, we additionally highlight its utility in the discovery of functional materials, in particular solar light absorbers.

Using the same V–Mn–Nb oxide library, we constructed a composition map of band gap energy using automated UV–vis spectroscopy¹⁶ and Tauc analysis.¹⁷ By correlating band gap energy with the phase concentration and shift parameter from the AgileFD–Gibbs solution, we have discovered new visible-band gap metal oxides and alloying-based tuning of the direct-allowed band gap energy. Such metal oxide semiconductors are critically missing in the quest to develop efficient solar fuel generators,¹⁸ and the combination of automated phase and band gap mapping described herein will significantly bolster the recent acceleration in the discovery of light absorbers for solar energy applications and beyond.¹⁹

■ ALGORITHM AND EXPERIMENTS

AgileFD Algorithm. In this article, we highlight the innovations of AgileFD, mainly focusing on its incorporation of physical constraints, which drastically differs from previous-reported approaches. For a comprehensive description of the mathematics of AgileFD, we refer the reader to ref.²⁰

Perhaps the most substantial advancement of AgileFD for the phase mapping problem is the efficient modeling of Peak Shifting, which we enable by applying a log-transformation to the XRD patterns. Each XRD pattern can be represented as the scattering intensity (I) as a function of q , the magnitude of the X-ray scattering vector. The contraction of a crystal lattice due to alloying can be considered as a scaling of the lattice parameter ($a \rightarrow \gamma a$), which corresponds to the inverse scaling of the scattering vector magnitude of each Bragg peak in the XRD pattern ($q_{\text{peak}} \rightarrow \gamma^{-1} q_{\text{peak}}$). For the present discussion, we consider lattice contraction ($\gamma < 1$) and note below that the method generalizes to lattice contraction and expansion. By performing a log-transformation of q , the peak shifting becomes ($\log q_{\text{peak}} \rightarrow \log \gamma^{-1} + \log q_{\text{peak}}$), and the additive shift can be incorporated in the matrix representation of XRD patterns through a row-shifting operation that we define below. The numerical values of q are not used in matrix factorization, but the q -spacing between data points is important when modeling shifting. A consistent model of Peak Shifting is enabled by choosing a constant interval ρ in the $\log q$ space so that the abscissa of the XRD patterns is a geometric series of q values

$$I_n(\rho^l q_{\text{min}}) \text{ where } n \in [0, N - 1], l \in [0, L - 1] \quad (2)$$

where the N XRD patterns correspond to N composition samples and are represented as the columns of the input data matrix A . By choosing a value of L and using the same range of

q values as the source data, these patterns are calculated by resampling the raw XRD patterns, and we typically choose L to be the number of data points in the raw XRD patterns, which is 2082 in the present work. A smaller value of L (larger value of ρ) can be chosen to accelerate calculations at the expense of q resolution. The logarithmic resampling results in decreasing q resolution as q increases, which may prompt the use of a larger value of L (smaller value of ρ) to retain the q resolution of the raw data. As noted below, the value of ρ also determines the peak shifting resolution.

In traditional NMF, the factorization produces K basis patterns W_k , where K is the number of phases used to describe the material system. Our innovative model to account for shifting is enabled by using $K \times M$ patterns in W that consist of M shifted versions of the basis pattern of each phase. Starting with the “non-shifted” version $W_{k,0}$, the additional $M - 1$ shifted versions are automatically generated in each optimization step by the row-shifting operation:

$$W_{k,m}(\rho^l q_{\text{min}}) = \begin{cases} W_{k,0}(\rho^{l-m} q_{\text{min}}) & \text{for } l \in [m, L - 1] \\ 0 & \text{for } l \in [0, m - 1] \end{cases}, m \in [1, M - 1] \quad (3)$$

Each shifted version corresponds to a lattice parameter contraction by a factor of $\gamma = \rho^{-m}$, which means that compared to the highest (or lowest) lattice parameter realized in the composition library, lattice contraction (or expansion) by $\rho^{-(M-1)}$ (or $\rho^{(M-1)}$) can be appropriately modeled by activating the corresponding shifted version of the basis pattern. That is, phase k is activated by $N \times M$ coefficients ($H_{n,k,m}$) in the matrix H such that the total activation of phase k for sample n is

$$h_{n,k} = \sum_{m=0}^{M-1} H_{n,k,m} \quad (4)$$

Since multiple terms in this sum can be nonzero, multiple shifted versions of each basis pattern can be used to create a pseudocontinuous model of peak shifting as long as the Bragg peaks are relatively wide ($\Delta q_{\text{peak}} > \rho \times q_{\text{peak}}$), which corresponds to using a sufficiently large value of L such that each peak spans over several data points. The representative shift parameter for a given sample and phase is calculated as a weighted mean

$$s_{n,k} = \langle \rho^m \rangle_{H_{n,k,m}} = \sum_{m=0}^{M-1} \rho^m H_{n,k,m} / h_{n,k} \quad (5)$$

which is the expectation value of ρ^m for each sample and for each basis pattern given the corresponding distribution of activation values.

While this construction of peak shifting is quite intuitive, the pattern shifting cannot be directly incorporated into NMF, which is one of the primary motivations for using cNMF as a starting point in developing AgileFD. Candidate solutions (typically starting with random seeding of H and W) are improved in AgileFD through scalable update rules where the loss function is translated into multiplicative matrix operations that enforce gradient descent. In ref.²⁰ we derive AgileFD's custom update rules for both H and W , which are derived from a loss function based on the generalized Kullback–Leibler divergence, an advancement that is critical for the phase mapping problem and is generally applicable for other source

separation problems. With these lightweight update rules, this incorporation of Peak Shifting is vastly more efficient than the previously proposed time warping techniques.^{7,14} The asymptotic time complexity for each solution update is $O(KLMN)$, or more precisely $\Theta(KLMN)$. No explicit bound exists for the number of updates needed until the completion of the gradient descent algorithm. The calculation time increases with decreasing convergence tolerance, requiring the user to define a convergence criterion that appropriately balances data reconstruction quality and convergence time, as discussed further below.

While we have developed and continue to develop a variety of update rules to customize the loss function and impose various constraints, we focus on two additional types of constraints in the present work, namely imposing additional expert constraints and enforcing Gibbs' Rule. Expert constraints comprise a variety of methods in which an expert can inject knowledge of the phase behavior, and AgileFD provides convenient mechanisms for encoding constraints through appropriate initialization of H and W . The multiplicative update rules have the convenient property that if any matrix value is zero in a candidate solution, it will remain zero throughout AgileFD convergence, which makes incorporation of certain constraints as simple as initializing specific matrix elements to zero. For example, to restrict the amount of shifting for a select basis pattern, the user can initialize the H values corresponding to the activation of some of its shifted copies to zero. To restrict the activation of a basis pattern to a select composition region, the user can make its activation zero for all other samples. A novel aspect of AgileFD is that imposing such constraints is substantially more computationally efficient than the (more expressive) incorporation of constraints in combinatorial reasoning algorithms such as CombiFD.¹¹ The most commonly desired expert constraint for W is to predefine the basis pattern for a phase that is known to exist in the composition library. The corresponding basis pattern can be incorporated in the initialization of W , and by not updating that portion of the matrix, the basis pattern is "frozen" in the AgileFD solution.

Gibbs' Rule is incorporated into AgileFD through an extension called AgileFD-Gibbs. In terms of the $h_{n,k}$ values, this rule corresponds to restricting the number of nonzero $h_{n,k}$ to be no greater than the number of components in the composition of sample n . The full Gibbs phase rule is actually more restrictive than the one implemented in AgileFD-Gibbs, but we find that restricting the number of coexisting phases according to the number of components is sufficient for practical use. The lack of enforcement of Gibbs' Rule has important consequences on the solution that cannot be ameliorated *ex post facto*. That is, a candidate solution can be "corrected" by zeroing the smallest $h_{n,k}$ values for each n until the solution satisfies Gibbs' Rule, but this guarantees that the corresponding basis patterns are not optimal and that the reconstructed data set may be "missing" Bragg peaks that are in the source data. More generally, this simple "correction" algorithm may not generate optimal solutions even after additional optimization of the basis patterns. To more elegantly and optimally enforce Gibbs' Rule, we combine AgileFD with a MIP algorithm that modifies a candidate solution. While W is fixed, the MIP algorithm generates a new activation matrix by finding the optimal $K \times M$ activation coefficients for each sample that adhere to Gibbs Rule, and the resulting candidate solution is then further optimized using AgileFD. The MIP

algorithm is applied to each sample *independently*, and since the complexity of the MIP is bound by the number of allowed phases in Gibbs' Rule, which is 3 in the present work and generally a small number, the most prominent scaling for the execution of the MIP step is its proportionality to N , the number of samples. It is worth noting the similarity between the utility of the MIP algorithm and the expert constraints noted above; the MIP-modified solution assigns the appropriate number of values in H to zero such that subsequent candidate solutions adhere to Gibbs' Rule.

With an AgileFD or AgileFD-Gibbs solution in hand, the collection of basis patterns $W_{k,0}$ and the composition maps of $h_{n,k}$ and $s_{n,k}$ provide an intuitive visualization of the phase behavior. Further quantification of phase concentrations and lattice parameters require identification of the crystal structure for each basis pattern, which was performed in the present work through search and match with the entries in the International Center for Diffraction Data (ICDD) database using DIF-FRAC.SUITE EVA software. Using relative intensity ratios and tabulated Bragg peaks from the ICDD entries (see [Supporting Information](#) for details), the relative total scattering intensity per mole of metal (ϑ_k) was calculated for each phase, enabling the relative molar activation of each phase in each sample to be calculated as

$$\text{rel}P_{n,k} = \frac{h_{n,k}}{\vartheta_k} \int W_{k,0} \quad (6)$$

where the integral over the basis pattern is analogous to the integration over the ICDD pattern. The k phase fractions for each sample are then provided through normalization of these relative molar activations:

$$P_{n,k} = \text{rel}P_{n,k} / \sum_{k=0}^{K-1} \text{rel}P_{n,k} \quad (7)$$

which is the solution to the phase map problem under the approximation that the total relative scattering factor of each phase matches that of the ICDD pattern and does not vary substantially with alloy composition within the phase.

To provide a more intuitive visualization of the weighted shift parameters, $s_{n,k}$, they are converted to lattice expansion parameters $\gamma_{n,k} = s_{n,k}^{-1}$. By determining the lattice expansion parameter $\gamma_{k,ICDD}$ that corresponds to the shift of $W_{k,0}$ that best matches the respective ICDD pattern, the relative shift of each phase in each sample compared to its ICDD entry is the ratio of $\gamma_{n,k}$ to $\gamma_{k,ICDD}$.

Library Synthesis. The continuous composition spread of V, Mn, and Nb was synthesized by reactive magnetron cosputtering in the presence of O_2 and Ar gas using elemental sources arranged symmetrically with respect to the 100 mm Si/SiO₂ substrate. The deposition proceeded for about 10 h with the RF power of the V, Mn, Nb source fixed at 150, 115, and 81 W, respectively. The spatial variation in deposition rate from each source resulted in a continuous thin film with composition gradient of the order of 0.5 at. % mm⁻¹. The variation in the deposition rates among different sources resulted in thickness variation across the library within a factor of 2 of the center thickness of 400 nm. The as-deposited oxide composition library was oxidized in air at 883 K for 1 h, producing the (V–Mn–Nb)O_x library for which oxygen composition is unknown and samples are represented by their cation composition.

Composition and Structure Measurements and Analysis. XRD data was acquired using a custom HiTp

setup incorporated into the bending-magnet beamline 1–5 of the Stanford Synchrotron Radiation Light source (SSRL) at SLAC National Accelerator Laboratory. A detailed description of the experiment is provided by Gregoire et al.,^{2a} and the characterization of the (V–Mn–Nb)O_x library employed a monochromated 13 keV source in reflection scattering geometry with an XRD image detector (Princeton Quad-RO 4320). Using a 1 mm² footprint, $N = 317$ samples were acquired on a square grid with 4.5 mm pitch and within a radius of 45 mm from the center of the 100 mm-diameter library. Diffraction images were processed into XRD patterns, $I_n(q)$, using WxDiff software²¹ and further processed with a custom background subtraction algorithm using cubic splines.

The XRF measurements were performed on an EDAX Orbis Micro-XRF system (EDAX Inc., Mahwah, NJ) with an X-ray beam approximately 1 mm in diameter. The V K, Mn K and Nb K XRF intensities were extracted from the Orbis software and converted to normalized V–Mn–Nb compositions using relative sensitivity factors calibrated at substrate center via a separate composition measurement. The calibration composition was measured on an Oxford Instruments X-Max 80 mm² energy dispersive X-ray spectroscopy (EDS; Oxford Instrument, Concord, MA) detector on a FEI Nova NanoSEM 450 (FEI, Hillsboro, OR). The absolute composition uncertainty for this EDS measurement is 10 at. %, and the XRF-determined relative compositions, which enable composition-property maps, have approximately 1 at. % resolution.^{4,22}

Optical Characterization. Optical characterization of the (V–Mn–Nb)O_x library was performed on a custom HiTp diffuse reflectance (DR) instrument built to analyze light absorber thin films on opaque substrates. The computer-automated experiment employed a 200 W Hg(Xe) lamp (Newport/Oriel Apex) and an integrating sphere (Ocean Optics ISP-50-8-R-GT) fiber-coupled to a spectrometer (Spectral Products, Inc. SM303), with further details provided by Mitrovic et al.¹⁶ DR spectra were acquired on a square grid of 1521 positions with a pitch of 2.032 mm, spot size of approximately 1 mm, and integration time of 0.025 s per spectrum, and typically three spectra were averaged to produce the DR spectrum for each sample. The absorption coefficient (α) was calculated up to a factor of a spectral scattering factor, which we approximate to be energy-independent, using the Kubelka–Munk radiative transfer model²³ which enabled calculation of the normalized Tauc property

$$TP = (\alpha h\nu)^{1/n} / \max((\alpha h\nu)^{1/n}) \quad (8)$$

where n is the Tauc exponent whose value is 1/2 for analysis of direct-allowed (DA) electronic transitions. The plots of DA Tauc property (TP^{DA}) versus photon energy ($h\nu$) were not only automatically generated but also automatically interpreted using a recently developed algorithm^{17b} that mimics expert judgment. The algorithm either estimates the DA band gap energy or determines that the band gap cannot be confidently estimated from the Tauc plot, which can occur when the band gap energy is beyond the range of the spectrometer, the sample is not sufficiently absorbing, or the direct-allowed band gap signature is convoluted by the presence of multiple phases with comparable contributions to the DR signal. We note that for discovery of photoabsorbers, indirect-allowed (IA) transitions are of primary interest only when the IA band gap is substantially lower than DA band gap. For the DA results shown here, no IA band gaps were identified at energies more than 0.3 eV below the direct-allowed gap.

The optical and XRF measurements were performed on the same set of library samples. The XRD measurements were performed on a coarser grid of library positions, and the composition for each XRD sample was calculated using linear interpolation in the Cartesian library position space. The results of XRD analysis ($\gamma_{n,k}$ and $P_{n,k}$) were interpolated to the optical and XRF samples through linear interpolation in the ternary composition space. The composition spread library and compositions of the samples characterized are shown in Figure 1.

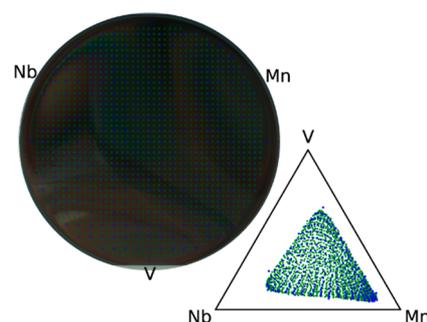


Figure 1. White light image of the (V–Mn–Nb)O_x composition library on 100 mm Si/SiO₂ substrate is shown with elemental labels indicating the orientation of the sputter deposition sources. The grid of library locations for XRD (blue) and both XRF and optical characterization (green) are shown along with an additional plot showing the XRF-determined compositions of these points.

RESULTS AND DISCUSSION

The phase behavior of the (V–Mn–Nb)O_x library was analyzed using AgileFD-Gibbs in unsupervised mode and then with additional expert constraints. Manual visualization of the XRD patterns revealed substantial Peak Shifting, prompting our use of $M = 10$ shifted copies of each basis pattern, which corresponds to a maximum lattice expansion of 1.0072 with respect to the basis pattern with the lowest lattice constant. Several values of K were attempted, revealing that the $K = 6$ solution appeared physically sound and that higher values of K neither substantially improved the data reconstruction nor produced basis patterns that were distinctive from those of the $K = 6$ solution. A more objective determination of appropriate values for parameters such as K and M is the subject of ongoing research and is beyond the scope of the present work.

Using the basis patterns from the $K = 6$ solution, we performed phase matching and identified the first 6 phases in Table 1. That is, upon fixing the values of K and M , the unsupervised AgileFD-Gibbs algorithm produced 6 basis patterns that match those of known phases despite substantial overlap of the basis patterns, which is a testament to the excellent source separation enabled by the update rules and Gibbs rule constraint. To demonstrate how the identification of 6 phases was enabled by the Peak Shifting feature of AgileFD, we show the AgileFD $M = 1$ and AgileFD-Gibbs $M = 10$ solutions in Figure 2. With $M = 1$, the dimension of the activation matrix H is equivalent to that of NMF, motivating our labeling of this solutions as “~NMF”, but the algorithms are not equivalent because of both the log transformation and gradient descent with custom update rules used in AgileFD. While some basis patterns and their phase distributions are similar between the AgileFD-Gibbs and ~NMF solutions, the lack of peak shift modeling has an important consequence in

Table 1. Primary Phases in the AgileFD-Gibbs Solutions for $K = 6$ and 8 Along with Their ICDD Entries^a

k , phase index	formula unit (crystal system)	ICDD entry number	relative total scattering intensity per mole of metal (θ_k)
0	Mn ₂ O ₃ (cubic)	01-071-0636	923.2
1	V _{2.38} Nb _{10.7} O _{32.7} (orthorhombic)	01-079-8393	2195.7
2	MnNb ₂ O ₆ (orthorhombic)	01-072-0484	2195.2
3	Mn ₃ V ₂ O ₈ (unknown)	00-039-0091	739.2 ²⁶
4	MnV ₂ O ₆ (monoclinic)	01-072-1837	911.8
5	Mn ₃ O ₄ (tetragonal)	01-080-0382	964.5
6	NbVO ₅ (orthorhombic)	00-046-0046	1706.2 ²⁷
7	V ₂ O ₅ (orthorhombic)	00-041-1426	678.0

^aICDD entries are overlaid in the basis pattern plots of Figures 2 and 3. The θ_k values used to calculate phase fractions are also shown.

the \sim NMF solution. Inspection of the $k = 5$ basis pattern in Figure 2 reveals that while the AgileFD-Gibbs solution is an excellent match to the ICDD pattern, the \sim NMF pattern is a very poor match. Instead of separating this ICDD phase, the data is better reconstructed in the \sim NMF solution by using 2 basis patterns ($k = 2$ and 5) to model the $k = 2$ basis pattern of the AgileFD-Gibbs solution. The $k = 2$ and 5 basis patterns in the \sim NMF solution are essentially shifted versions of each other and activated in complementary portions of the composition space. The \sim NMF solution fails to identify the $k = 5$ phase from the AgileFD-Gibbs solution as a direct consequence of the inability to model Peak Shifting.

Additional subtle but important difference exist between the \sim NMF and AgileFD-Gibbs solution, and perhaps the illustrative example for the present discussion is that the

AgileFD-Gibbs solution adheres to the connectivity requirement much more closely than the \sim NMF solution even though this constraint was not imposed. The enforcement of the Gibbs rule has important implications for optimizing not only the phase distributions but also the basis patterns. It is important to note that the quality of the phase map solution lies in its explanation of physically meaningful phase behavior, which requires (at a minimum) that the basis patterns represent solid state phases with plausible phase distributions in composition space. Considering the reconstruction error for each sample to be the sum of the absolute residuals (E), then with the total intensity in the XRD pattern (T), the fractional reconstruction accuracy can be calculated as $1 - E/T$. Averaged over the samples, the \sim NMF and AgileFD-Gibbs solutions have 73% and 79% accuracy, respectively, but the improvements in the AgileFD-Gibbs solution are not sufficiently represented by this increase in reconstruction accuracy. Indeed a challenging aspect of the phase demixing problem is that a multitude of solutions provide similar reconstruction accuracy, requiring the incorporation of physical constraints to yield a solution that is both accurate and physically meaningful.

The AgileFD-Gibbs solution in Figure 2 reveals that the lattice constant for several phases varies systematically through the composition space, a strong indicator of alloying. Upon careful inspection of the AgileFD-Gibbs solution we also noticed some opportunities to further improve the phase map via expert constraints in AgileFD-Gibbs. We proceeded to manually inspect the compositions closest to the end-member binary oxides. For Mn oxide the AgileFD-Gibbs solution produced 2 different Mn oxide phases ($k = 0$ and 5). We have commonly observed the coexistence of these phases in Mn-containing libraries of metal oxides due to the lack of a strong thermodynamic differentiation under the library synthesis conditions (Gibbs free energies of formation of -359 and -356 kJ mol⁻¹ per Mn atom for Mn₂O₃ and Mn₃O₄, respectively). Since the library compositions contain substantial

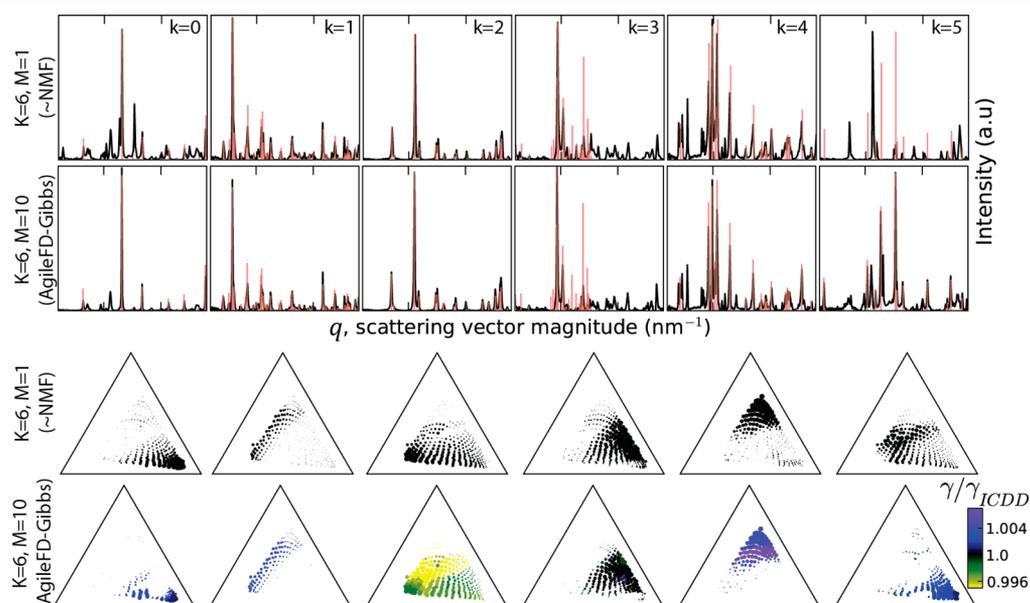


Figure 2. Six-phase solutions for the (V–Mn–Nb)O_x library with algorithm parameters noted in the left-hand labels. The \sim NMF solution is the AgileFD solution with the peak shifting removed by setting $M = 1$. The basis patterns are plotted along with the ICDD patterns listed in Table 1. The map of each phase is shown as a composition plot where the point size represents the phase fraction $P_{n,k}$ and the color represents the relative lattice constant compared to the respective basis pattern, which is aligned to the best-match with the ICDD pattern.

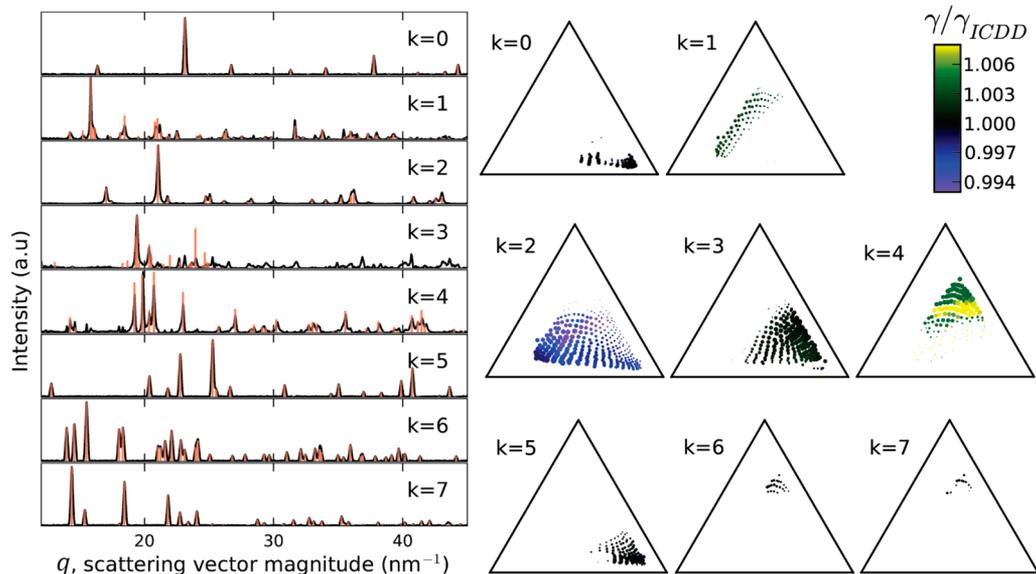


Figure 3. AgileFD-Gibbs solution with $K = 8$, $M = 10$, and expert constraints applied to $k = 0, 5, 6, 7$. The map of each phase is shown as a composition plot where the point size represents the phase fraction $P_{n,k}$ and the color represents the relative lattice constant, as described in Figure 2.

amounts of Nb and V that may alloy into these phases, the alloyed metal oxides may indeed coexist in thermodynamic equilibrium in this ternary composition space. Through manual inspection of XRD patterns containing these phases, we confirmed that both phases are present and in fact they often coexist. It is worth noting that without prior knowledge of the XRD patterns of pure phases, any source separation algorithm cannot robustly produce phase-pure basis patterns for phases that always coexist. For the Mn oxides, AgileFD-Gibbs provides sufficient separation of the phases that they were identified in ICDD search and match, and we enforced complete separation by seeding W with simulated basis patterns obtained by convoluting the ICDD patterns with a Gaussian peak shape ($\sigma = 0.13 \text{ nm}^{-1}$). Inspection of the most Nb-rich samples revealed that, as suggested by Figure 2, the library samples are far enough from the Nb end-member composition that no Nb oxides are observed. Inspection of the V-rich composition revealed the presence of 2 minor phases, V_2O_5 and NbVO_5 , whose relatively weak intensity and existence in only a small fraction of the composition samples gives their signals little contribution to the system-wide loss function. While AgileFD-Gibbs may factor these phases with an extended data set that includes compositions closer to the stoichiometries of these phases, we amended this shortcoming of the source data by introducing and freezing a new basis pattern for both V_2O_5 and NbVO_5 . Since these phases were only observed in V-rich compositions, we additionally seeded the corresponding values of H with zeros for all samples with V concentration less than 0.45. The expanded $K = 8$ solution with expert constraints on four phases was generated using AgileFD-Gibbs, producing the solution shown in Figure 3.

The addition of 2 phases and expert constraints enabled more accurate modeling of every phase, leading to slight changes (compared to the AgileFD-Gibbs solution of Figure 2) in the basis patterns and phase maps even for the phases that were not constrained. The agreement between $k = 1, 2, 4$ and the respective basis patterns is remarkable, and the agreement for $k = 3$ is quite good but not perfect. Upon manual inspection of samples with high activation of this phase, we believe this basis pattern accurately models the complex phase behavior in

this composition region. It appears that the $\text{Mn}_3\text{V}_2\text{O}_8$ coexists with minor fractions of 2 polytypes of $\text{Mn}_2\text{V}_2\text{O}_7$. We previously reported²⁴ on the coexistence of these polytypes in thin film $(\text{V-Mn})\text{O}_x$ libraries and did not pursue further disambiguation in the phase map solution since all 3 of these manganese vanadates only appear as minor phases since the composition samples are all beyond ~ 10 at. % of the V–Mn binary line.

With the phase map of Figure 3 in hand, we turn to the optical characterization data to investigate if the composition library contains any metal oxide phases of interest for solar light absorption. While the automated Tauc analysis algorithm did not identify a DA transition for some samples, DA band gap energies (E_g^{DA}) were estimated for 1329 samples and are mapped in composition space in Figure 4. While comparisons

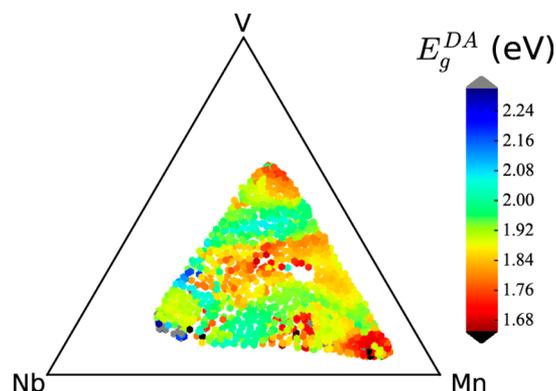


Figure 4. Composition map of the DA band gap energy from automated Tauc analysis is shown for 1329 composition samples.

of the band gap data and phase map reveal a number of interesting structure–property relationships, for the present purposes we provide a detailed interpretation of MnV_2O_6 , corresponding to $k = 4$ in Figure 3. It is worth noting that for mixed-phase samples, the majority phase does not necessarily provide the majority contribution to the Tauc plot due to possible differences in the optical absorption strength and band gap energies of different phases. To infer which composition

samples exhibit a band gap representative of MnV_2O_6 , we show the band gap energy as a function of MnV_2O_6 phase fraction in Figure 5, in which the sample points are colored according to

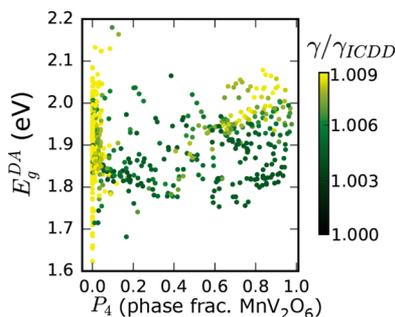


Figure 5. DA band gap energy is plotted against the $k = 4$ phase fraction from Figure 3, demonstrating that when the phase fraction of MnV_2O_6 is above 0.8, the band gap value systematically varies with the relative lattice parameter.

the relative lattice constant from the $k = 4$ phase map of Figure 3. With low phase fraction, the band gap energy is most likely due to a different phase, so inspecting the convergence of the band gap energy as the phase fraction increases provides an indication of the phase fraction at which MnV_2O_6 is providing the major contribution to the band gap measurement. Figure 5 reveals that samples with phase fraction in excess of 0.8 exhibit a band gap in the 1.8–2.08 range, and that the variation within this range is correlated with the lattice parameter.

To further investigate this relationship between band gap and lattice parameter in MnV_2O_6 , we extract the samples with phase purity in excess of 0.8 and visualize in Figure 6 both the original XRD patterns and the relative lattice parameter from AgileFD-Gibbs. This visualization provides a more detailed understanding of the phase behavior in this composition region. At the lowest V concentration (composition $\text{V}_{0.19}\text{Mn}_{0.38}\text{Nb}_{0.43}\text{O}_x$), the lattice parameter is approximately 1.009 times larger than the ICDD entry and as the V concentration is increased to 0.5, the relative lattice constant systematically lowers and trends toward 1 as the cation composition gets closer to the $\text{Mn}_{0.33}\text{V}_{0.67}$ composition of stoichiometric MnV_2O_6 . The composition library does not extend to this composition, and in the composition region noted in Figure 6, there is a phase boundary at V concentration of 0.5 where the alloying of MnV_2O_6 does not continue into more V-rich compositions and this phase becomes mixed with NbVO_5 . The highlighted composition region in Figure 6 contains phase-pure MnV_2O_6 (within the detectability limit of other phases) and exhibits alloying-based peak shifting with respect to both Mn and Nb, suggesting that the Mn concentration varies by ± 4 at. % from the phase stoichiometry and that Nb is soluble up to approximately 25 at. %. For this noncubic crystal structure, the lattice constant determined by AgileFD-Gibbs is an “average” lattice constant but not necessarily the mean of the three lattice parameters. Figure 6a shows that most peaks shift uniformly with some indication of nonisotropic lattice expansion that we do not model in the present work and may provide insights into the elemental site substitutions of this oxide alloy.

To investigate the impact of alloying on band gap energy, we consider the samples in the composition region highlighted in Figure 6. As shown in Figure 7, the band gap energy of these samples spans a nearly 0.2 eV range and systematically increases

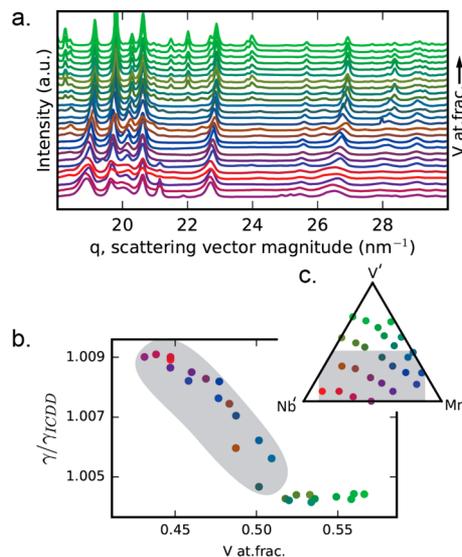


Figure 6. (a) For the 26 samples with phase fraction of MnV_2O_6 in excess of 0.8, the series of XRD patterns (in the q -range with primary ICDD peaks) are arranged according to the V concentration. The top 10 patterns show a small amount of NbVO_5 , and the bottom 16 patterns exhibit systematic shifting of the MnV_2O_6 peaks with respect to V concentration. (b) The relative lattice constant for the 26 samples is shown and the 16 samples with no NbVO_5 are indicated by a gray region. (c) The 26 samples are shown in a composition plot with end-members $\text{V}' = \text{V}_{0.61}\text{Mn}_{0.29}\text{Nb}_{0.09}\text{O}_x$, $\text{Mn}' = \text{V}_{0.43}\text{Mn}_{0.47}\text{Nb}_{0.1}\text{O}_x$ and $\text{Nb}' = \text{V}_{0.43}\text{Mn}_{0.29}\text{Nb}_{0.28}\text{O}_x$. The composition gray composition region contains the same 16 samples as that in b. The data for each sample is colored according to its composition in all 3 plots.

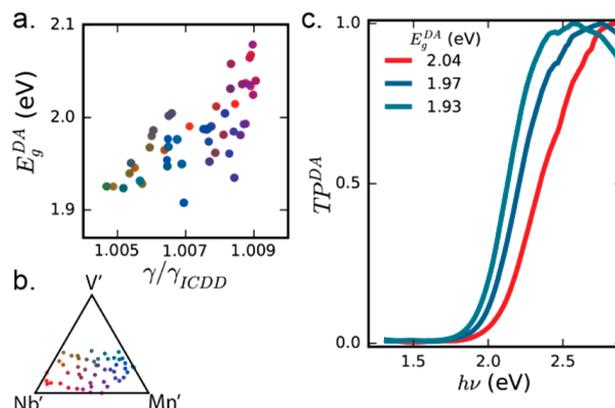


Figure 7. (a) Band gap values from the composition region containing high phase purity MnV_2O_6 (see Figure 6) are plotted against the relative lattice parameter, demonstrating alloying-based tuning of the band gap energy. (b) Sample compositions using the same composition-color scale revealing that among these samples, the highest band gap energy and lattice parameters are found with the most Nb-rich compositions. (c) Representative Tauc plots for 3 samples with line colors matching the samples' colors in b. The band gap energies produced by the automated Tauc algorithm are listed in the legend.

with increasing relative lattice parameter, which is highly correlated with composition as shown in Figure 6. We note that while the band gaps of these oxide alloys have not been previously reported, the trend in Figure 7 is commensurate with reported direct-allowed band gap value of 1.8–1.95 eV for MnV_2O_6 (where the relative lattice parameter is 1).²⁵ This alloying-induced variation in the band gap energy is well-

studied in semiconductors based on main-group elements but is rarely observed in metal oxides, particularly in the 1.8–2.1 eV range that is of primary interest for photoanodes in solar hydrogen generators and other applications.¹⁸ While substantial additional research is required to fully understand the alloying and band gap tuning in V–Mn–Nb oxides, the results presented in Figures 2–7 highlight the excellent performance of the AgileFD suite of algorithms and their enabling ability to conduct materials science research in high order composition spaces.

Since the matrix factorization in AgileFD proceeds via gradient descent from randomized initial matrices (other than the matrix elements set according to expert constraints), different randomizations may require different calculation times to reach the convergence criterion and may yield different solutions. It is important to ensure that a solution is “stable” by comparing the results from several random initializations, as was done for all solutions presented in the present work. Solutions equivalent (within negligible variations related to random initialization) to that of Figure 3 can be generated using single or multithreaded processing with convergence times of approximately 1200 and 370 s, respectively (Intel x5690 3.46 GHz processor), which in the multithreaded case is approximately equally distributed among the initial matrix factorization, the MIP solution for applying Gibbs’ Rule, and the subsequent matrix factorization. For this publication-quality solution, a stringent convergence criterion was used wherein the matrix updates proceeded until the change in the loss function dropped below 0.008%. It is worth noting that for initial exploration of data, a typical convergence threshold of 0.05% can be used, in which case the AgileFD matrix factorization completes in approximately 50 s for this sizable data set with KLMN = 52.8 million matrix elements.

In closing, we note some limitations of phase mapping techniques and algorithms that may be addressed in future research. It is worth noting that sputter deposition is a nonthermal synthesis technique, and even with postdeposition annealing of a library such as that described here, deviations from thermodynamic equilibrium may persist. As noted above, the coexistence of Mn₂O₃ and Mn₃O₄ may be an example of such nonequilibrium phase behavior, requiring that constraints based on Gibbs’ phase rule must be applied with care. To address this potential issue we routinely generate an AgileFD–Gibbs solution with an additional allowed phase (4 coexisting phases for ternary compositions) for comparison. Concerning the modeling of alloying-based peak shifting, AgileFD algorithm is somewhat limited as it is intended to track peak shifting resulting from isotropic lattice distortions. More complex alloying-based distortions may result in different Bragg peaks shifting by different amounts and even in different directions. Such peaks shifting can only be modeled in AgileFD by splitting the pattern of a phase into subpatterns representing sets of peaks that shift together, and through recognition that these “sub” basis patterns are activated over the same composition ranges, the full pattern can be reconstructed in post processing. This approach to modeling peak shifting is not commensurate with Gibbs’ phase rule and is not a sufficiently general model for alloying in noncubic systems, motivating further algorithm development. Finally, we noted above that adherence to the connectivity constraint is a general requirement for meaningful solutions, which is not enforced in AgileFD. In the V–Mn–Nb oxide system described here, we demonstrated that the AgileFD–Gibbs algorithm sufficiently

addresses the connectivity shortcomings of the ~NMF algorithm, but indeed a connectivity constraint is a desirable extension of AgileFD that will be addressed in future work. The source data can be found in the [Supporting Information](#) and at <http://www.udiscover.it/resources/data/>. The source code will be hosted at <http://www.udiscover.it/resources/software/>.

SUMMARY

We introduce AgileFD, the first scalable phase mapping algorithm for combinatorial XRD data that models alloying-based peak shifting and imposes Gibbs’ phase rule. The importance of encoding these properties of solid state phase diagrams into a source separation algorithm is discussed using the (V–Mn–Nb)O_x system where unsupervised mapping of six phases is directly enabled by these capabilities of AgileFD. Several of the 317 XRD patterns from the composition library contain small signals from minor phases that are not amenable to unsupervised factorization but can be modeled by AgileFD through straightforward injection of expert knowledge. The resulting 8-phase solution reveals alloying in ternary oxide phases such as MnV₂O₆. By combining the AgileFD solution with band gap energies obtained from automated Tauc analysis of high throughput UV–vis spectroscopy data, we identify band gap tuning of nearly 0.2 eV as a function of lattice parameter and V composition in the energy range of interest for solar applications. The identification of this family of promising solar light absorbers was enabled through a multidisciplinary effort in which materials-motivated advancements of computer science techniques produced powerful new algorithms for accelerating materials discovery.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: [10.1021/acscombsci.6b00153](https://doi.org/10.1021/acscombsci.6b00153).

XRD and composition data for the V–Mn–Nb oxide system and additional source data (ZIP)

AUTHOR INFORMATION

Corresponding Authors

*E-mail: gomes@cs.cornell.edu.

*E-mail: gregoire@caltech.edu.

ORCID

John M. Gregoire: [0000-0002-2863-5265](https://orcid.org/0000-0002-2863-5265)

Present Address

¹Ronan Le Bras: Allen Institute for Artificial Intelligence, Seattle, Washington 98103, United States

Notes

The authors declare no competing financial interest.

ACKNOWLEDGMENTS

The experimental work was performed in the Joint Center for Artificial Photosynthesis (JCAP), a DOE Energy Innovation Hub supported through the Office of Science of the U.S. Department of Energy under Award No. DE-SC0004993. The algorithm development is supported by NSF awards CCF-1522054 and CNS-0832782 (Expeditions), CNS-1059284 (Infrastructure), and IIS-1344201 (INSPIRE); and ARO award W911-NF-14-1-0498. Use of the Stanford Synchrotron Radiation Lightsource, SLAC National Accelerator Laboratory,

is supported by the U.S. Department of Energy, Office of Science, Office of Basic Energy Sciences under Contract No. DE-AC02-76SF00515. The authors thank Apurva Mehta and Douglas G. Van Campen for assistance with collection of synchrotron XRD data.

REFERENCES

- (1) Green, M. L.; Takeuchi, I.; Hatrick-Simpers, J. R. Applications of high throughput (combinatorial) methodologies to electronic, magnetic, optical, and energy-related materials. *J. Appl. Phys.* **2013**, *113* (23), 231101–231101.
- (2) (a) Gregoire, J. M.; Van Campen, D. G.; Miller, C. E.; Jones, R.; Suram, S. K.; Mehta, A. High Throughput Synchrotron X-ray Diffraction for Combinatorial Phase Mapping. *J. Synchrotron Radiat.* **2014**, *21* (6), 1262–1268. (b) Gregoire, J. M.; Dale, D.; Kazimirov, A.; DiSalvo, F. J.; van Dover, R. B. High energy x-ray diffraction/x-ray fluorescence spectroscopy for high-throughput analysis of composition spread thin films. *Rev. Sci. Instrum.* **2009**, *80* (12), 123905–123905.
- (3) Barr, G.; Dong, W.; Gilmore, C. J. PolySNAP: a computer program for analysing high-throughput powder diffraction data. *J. Appl. Crystallogr.* **2004**, *37* (4), 658–664.
- (4) Hatrick-Simpers, J. R.; Gregoire, J. M.; Kusne, A. G. Perspective: Composition–structure–property mapping in high-throughput experiments: Turning data into knowledge. *APL Mater.* **2016**, *4* (5), 053211.
- (5) (a) Bunn, J. K.; Han, S.; Zhang, Y.; Tong, Y.; Hu, J.; Hatrick-Simpers, J. R. Generalized machine learning technique for automatic phase attribution in time variant high-throughput experimental studies. *J. Mater. Res.* **2015**, *30* (7), 879–889. (b) Bunn, J. K.; Hu, J.; Hatrick-Simpers, J. R. Semi-Supervised Approach to Phase Identification from Combinatorial Sample Diffraction Patterns. *JOM* **2016**, *68* (8), 2116–2125.
- (6) Long, C. J.; Hatrick-Simpers, J.; Murakami, M.; Srivastava, R. C.; Takeuchi, I.; Karen, V. L.; Li, X. Rapid structural mapping of ternary metallic alloy systems using the combinatorial approach and cluster analysis. *Rev. Sci. Instrum.* **2007**, *78* (7), 072217.
- (7) Lebras, R.; Damoulas, T.; Gregoire, J. M.; Sabharwal, A.; Gomes, C. P.; van Dover, R. B. Constraint Reasoning and Kernel Clustering for Pattern Decomposition With Scaling. *Proc. 17th Intl. Conf. on Principles and Practice of Constraint Programming* **2011**, 6876, 508–522.
- (8) Kusne, A. G.; Gao, T.; Mehta, A.; Ke, L.; Nguyen, M. C.; Ho, K.-M.; Antropov, V.; Wang, C.-Z.; Kramer, M. J.; Long, C.; Takeuchi, I. On-the-fly machine-learning for high-throughput experiments: search for rare-earth-free permanent magnets. *Sci. Rep.* **2014**, *4*, 6367.
- (9) Ermon, S.; Le Bras, R.; Gomes, C. P.; Selman, B.; van Dover, R. B. SMT-Aided Combinatorial Materials Discovery. *Proceedings of the 15th International Conference on Theory and Applications of Satisfiability Testing (SAT)* **2012**, 7317, 172–185.
- (10) Long, C. J.; Bunker, D.; Li, X.; Karen, V. L.; Takeuchi, I. Rapid identification of structural phases in combinatorial thin-film libraries using x-ray diffraction and non-negative matrix factorization. *Rev. Sci. Instrum.* **2009**, *80* (10), 103902–103902.
- (11) Ermon, S.; Le Bras, R.; Suram, S. K.; Gregoire, J. M.; Gomes, C. P.; Selman, B.; Van Dover, R. B. Pattern Decomposition with Complex Combinatorial Constraints: Application to Materials Discovery. *AAAI'15 Proceedings of the Twenty-Ninth Conference on Artificial Intelligence (AAAI-15)* **2015**, 636–643.
- (12) Massoudifar, P.; Rangarajan, A.; Zare, A.; Gader, P. An integrated graph cuts segmentation and piece-wise convex unmixing approach for hyperspectral imaging. Presented at the 6th Workshop on Hyperspectral Image and Signal Processing (WHISPERS): Evolution in Remote Sensing, 24–27 June 2014, Lausanne, Switzerland; http://www.ieee-whispers.com/images/techprograms/whispers_2014_technical_program.pdf.
- (13) Kusne, A. G.; Keller, D.; Anderson, A.; Zaban, A.; Takeuchi, I. High-throughput determination of structural phase diagram and constituent phases using GRENDL. *Nanotechnology* **2015**, *26* (44), 444002.
- (14) Baumes, L. A.; Moliner, M.; Nicoloyannis, N.; Corma, A. A reliable methodology for high throughput identification of a mixture of crystallographic phases from powder X-ray diffraction data. *CrystEngComm* **2008**, *10* (10), 1321–1321.
- (15) Smaragdis, P. Non-negative Matrix Factor Deconvolution; Extraction of Multiple Sound Sources from Monophonic Inputs. *International Conference on Independent Component Analysis and Signal Separation* **2004**, 3195, 494–499.
- (16) Mitrovic, S.; Cornell, E. W.; Marcin, M. R.; Jones, R. J.; Newhouse, P. F.; Suram, S. K.; Jin, J.; Gregoire, J. M. High-throughput on-the-fly scanning ultraviolet-visible dual-sphere spectrometer. *Rev. Sci. Instrum.* **2015**, *86* (1), 013904.
- (17) (a) Suram, S. K.; Newhouse, P. F.; Zhou, L.; Van Campen, D. G.; Mehta, A.; Gregoire, J. M. High Throughput Light Absorber Discovery, Part 2: Establishing Structure-Band Gap Energy Relationships. *ACS Comb. Sci.* **2016**, *18*, 682. (b) Suram, S. K.; Newhouse, P. F.; Gregoire, J. M. High Throughput Light Absorber Discovery, Part 1: An Algorithm for Automated Tauc Analysis. *ACS Comb. Sci.* **2016**, *18*, 673.
- (18) Hu, S.; Xiang, C.; Haussener, S.; Berger, A. D.; Lewis, N. S. An analysis of the optimal band gaps of light absorbers in integrated tandem photoelectrochemical water-splitting systems. *Energy Environ. Sci.* **2013**, *6* (10), 2984–2984.
- (19) (a) Zhou, L.; Yan, Q.; Shinde, A.; Guevarra, D.; Newhouse, P. F.; Becerra-Stasiewicz, N.; Chatman, S. M.; Haber, J. A.; Neaton, J. B.; Gregoire, J. M. High Throughput Discovery of Solar Fuels Photoanodes in the CuO–V₂O₅ System. *Adv. Energy Mater.* **2015**, *5*, 1500968. (b) Thienhaus, S.; Naujoks, D.; Pftzing-Micklich, J.; Konig, D.; Ludwig, A. Rapid identification of areas of interest in thin film materials libraries by combining electrical, optical, X-ray diffraction, and mechanical high-throughput measurements: a case study for the system Ni–Al. *ACS Comb. Sci.* **2014**, *16* (12), 686–94. (c) Subramanian, A.; Perkins, J. D.; O'Hayre, R. P.; Ginley, D. S.; Lany, S.; Zakutayev, A. Non-equilibrium synthesis, structure, and optoelectronic properties of Cu_{2–2x}Zn_xO alloys. *J. Mater. Sci.* **2015**, *50* (3), 1350–1357.
- (20) Xue, Y.; Bai, J.; Le Bras, R.; Rappazzo, B.; Bernstein, R.; Bjorck, J.; Longpre, L.; Suram, S. K.; van Dover, R. B.; Gregoire, J.; Gomes, C. P. Phase-Mapper: An AI Platform to Accelerate High Throughput Materials Discovery. 2016, arXiv:1610.00689. arXiv.org e-Print archive. <http://adsabs.harvard.edu/abs/2016arXiv161000689X>.
- (21) Mannsfeld, S. C.; Tang, M. L.; Bao, Z. Thin film structure of triisopropylsilyl ethynyl-functionalized pentacene and tetraceno[2,3-b]thiophene from grazing incidence X-ray diffraction. *Adv. Mater.* **2011**, *23* (1), 127–31.
- (22) Le Bras, R.; Bernstein, R.; Gregoire, J. M.; Suram, S. K.; Gomes, C. P.; Selman, B.; van Dover, R. B. A Computational Challenge Problem in Materials Discovery: Synthetic Problem Generator and Real-World Datasets. Presented at the Twenty-Eighth Conference on Artificial Intelligence (AAAI-14), Québec City, Canada, 2014.
- (23) Kubelka, P. New Contributions to the Optics of Intensely Light-Scattering Materials. Part I. *J. Opt. Soc. Am.* **1948**, *38* (5), 448–457.
- (24) Yan, Q.; Li, G.; Newhouse, P. F.; Yu, J.; Persson, K. A.; Gregoire, J. M.; Neaton, J. B. Mn₂V₂O₇: An Earth Abundant Light Absorber for Solar Water Splitting. *Adv. En Mater.* **2015**, *5* (8), 1401840.
- (25) Zhang, W.; Shi, L.; Tang, K.; Liu, Z. Synthesis, surface group modification of 3D MnV₂O₆ nanostructures and adsorption effect on Rhodamine B. *Mater. Res. Bull.* **2012**, *47* (7), 1725–1733.
- (26) This ICDD entry does not have relative scattering information. The average value from Mn₂V₂O₇ polytypes (01-073-6361, 01-089-0483) is used as a proxy.
- (27) This ICDD entry does not have relative scattering information. The average value from V₂O₅ (00-041-1426) and Nb₂O₅ (04-007-2424) is used as a proxy.